

The workshop on partial differential equations on the sphere 2019  
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# Application of spherical helix nodes to a shallow-water model using radial basis functions



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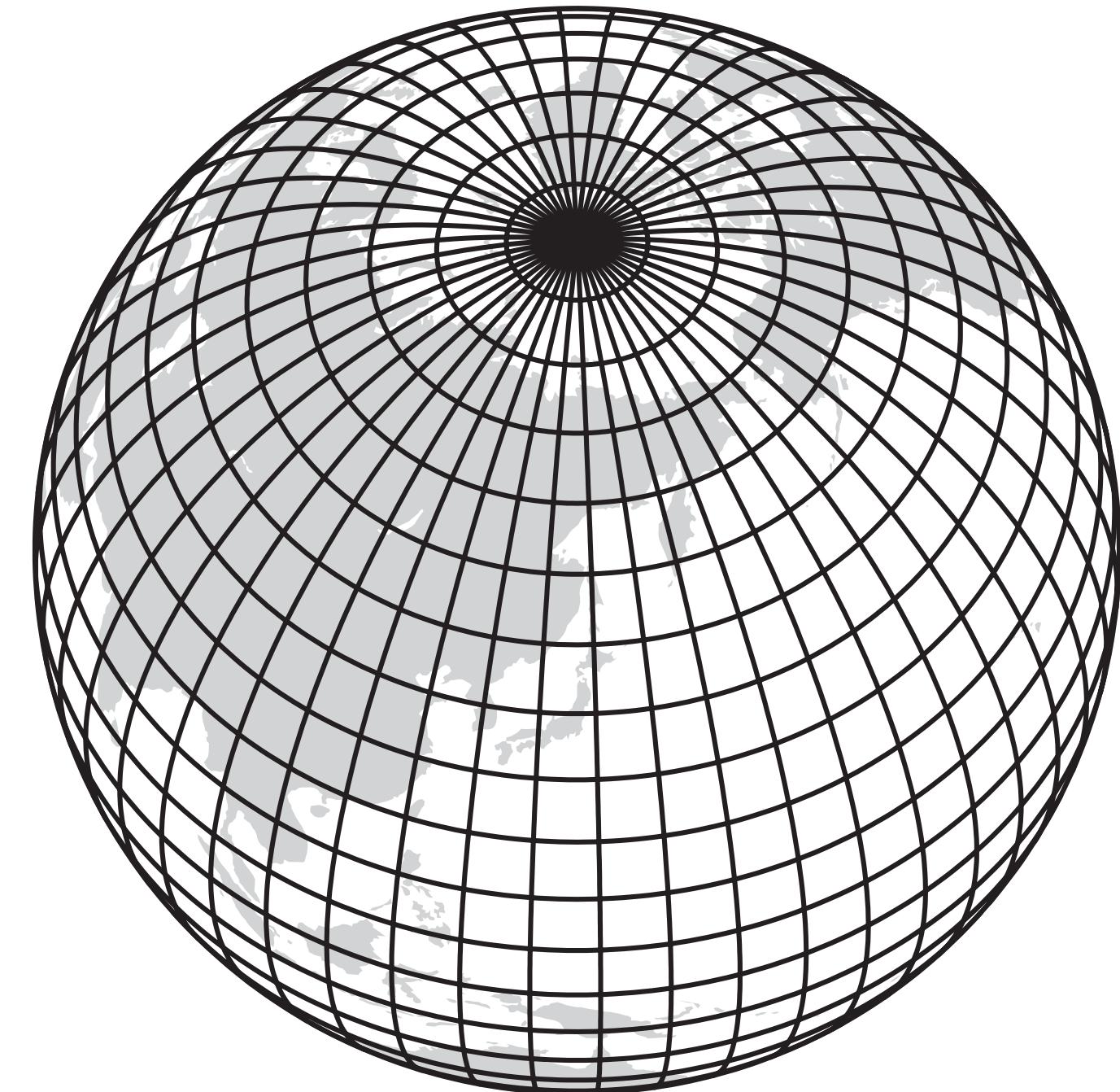
Supported by JSPS KAKENHI Grant Number 15K13417 and  
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(Elucidation of the Birth of Exoplanets [Second Earth] and  
the Environmental Variations of Planets in the Solar System)



# Pole problem

- Severe limitation of time-step due to the small zonal interval near the pole with lat-lon grid, calling for quasi-uniform distribution
- However, no obvious methods known to distribute many points uniformly on the sphere
- One of 18 mathematical problems left for 21C (Smale 1998)

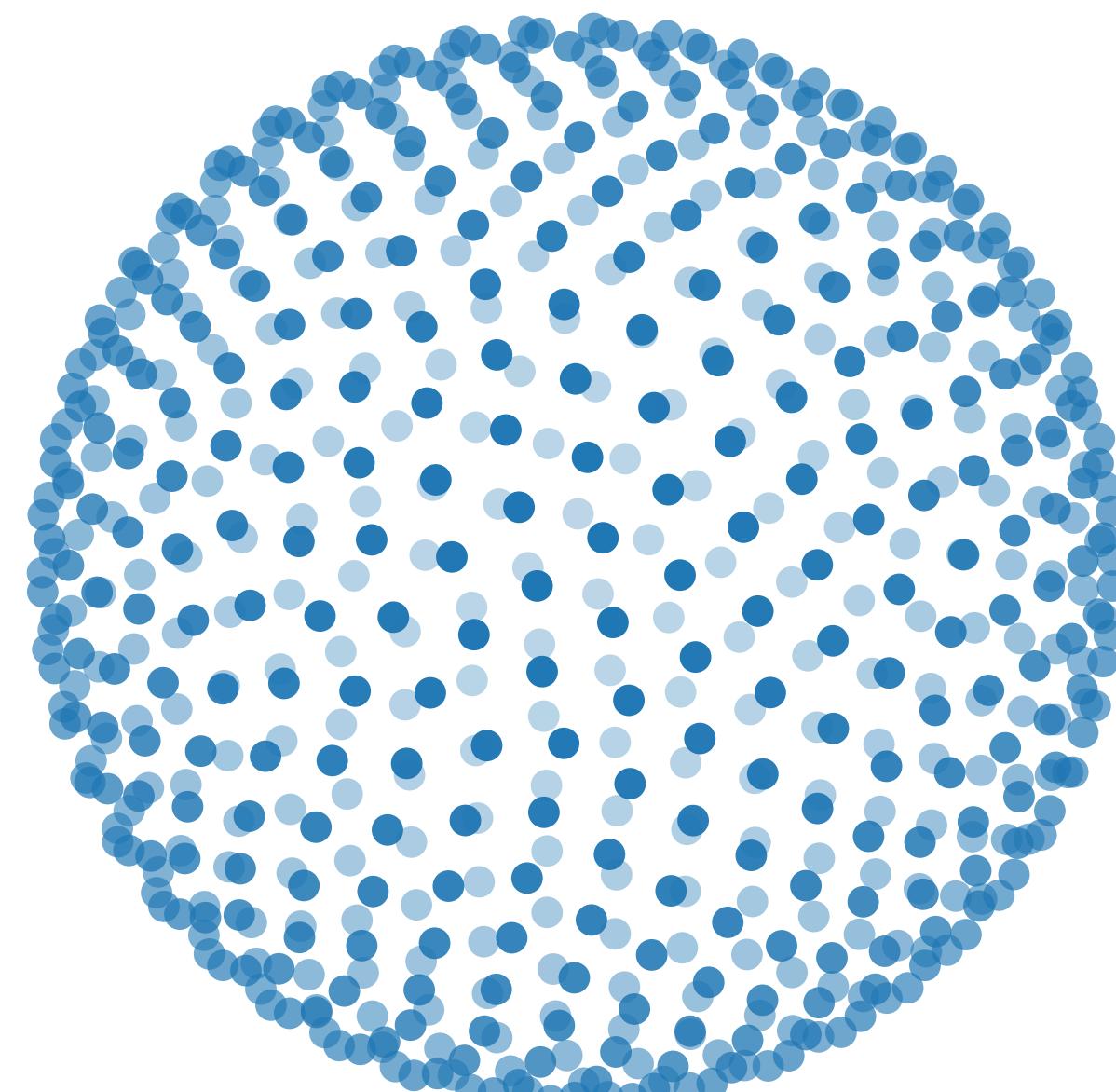
LATITUDE-LONGITUDE GRID



Williamson 2007

# Quasi-uniform nodes on the sphere

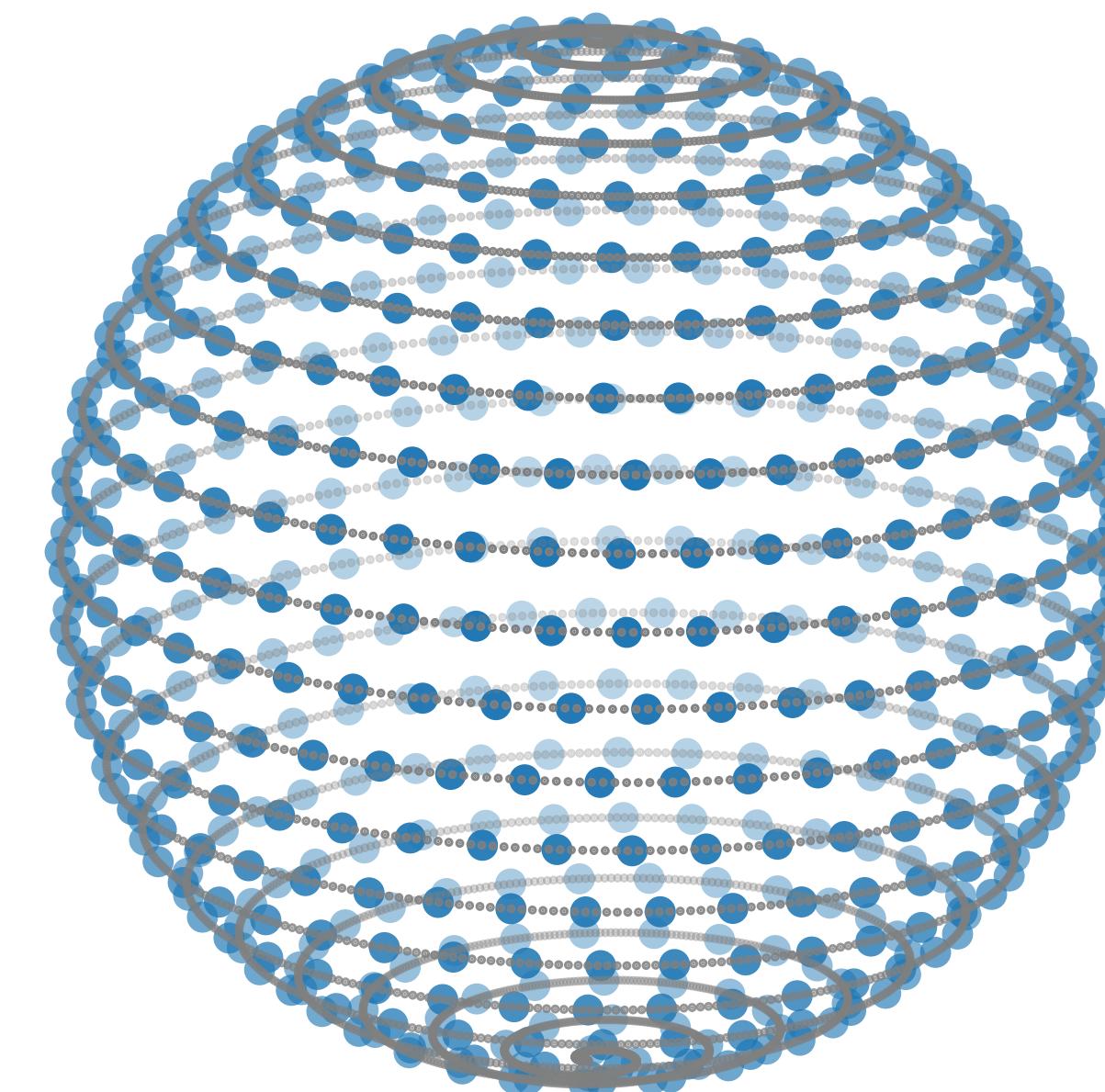
ME: Minimum Energy



[https://github.com/  
gradywright/spherepts](https://github.com/gradywright/spherepts)

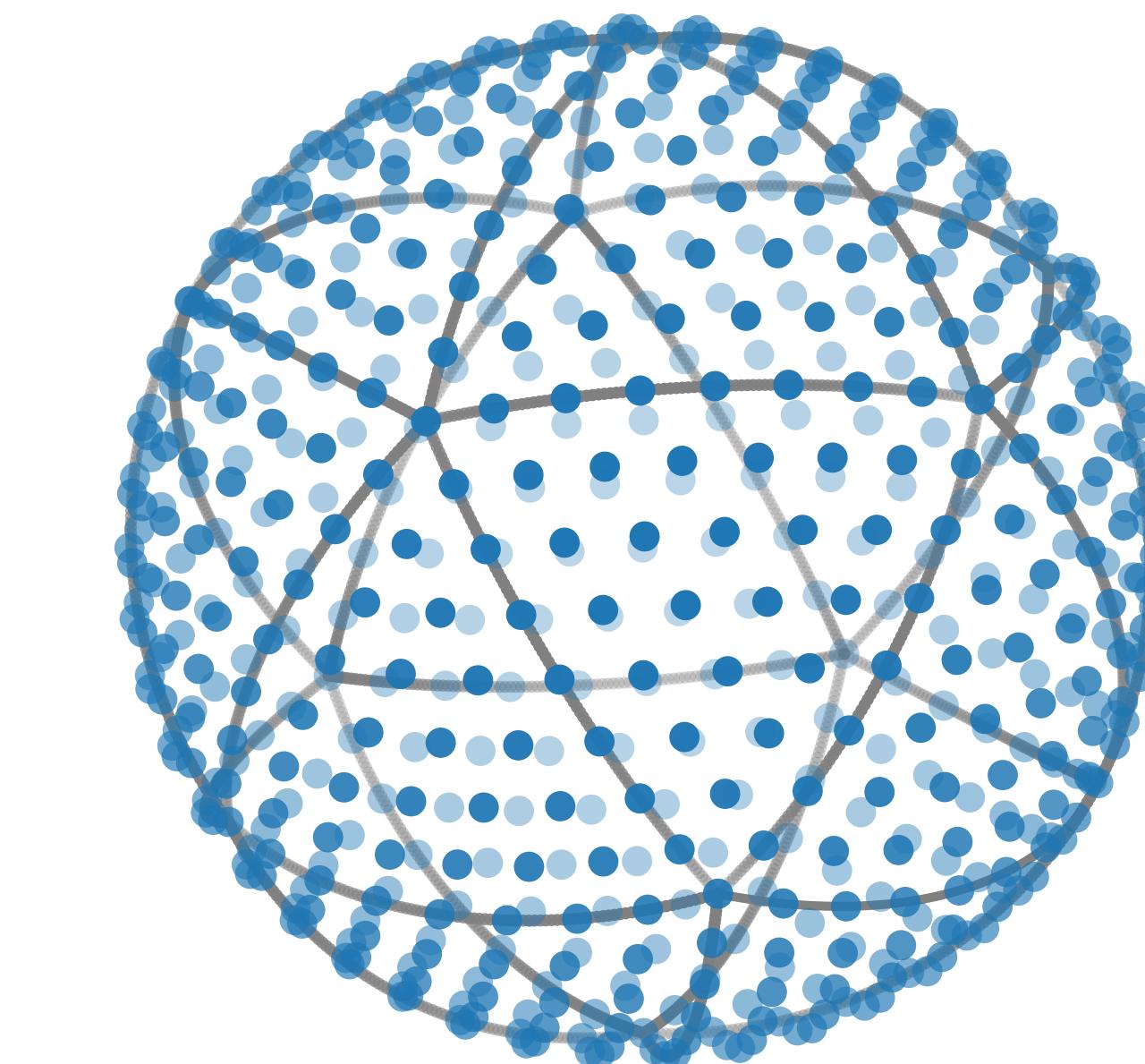
SH: Spherical Helix

$$\lambda = m\theta \mod 2\pi$$



Bauer 2000

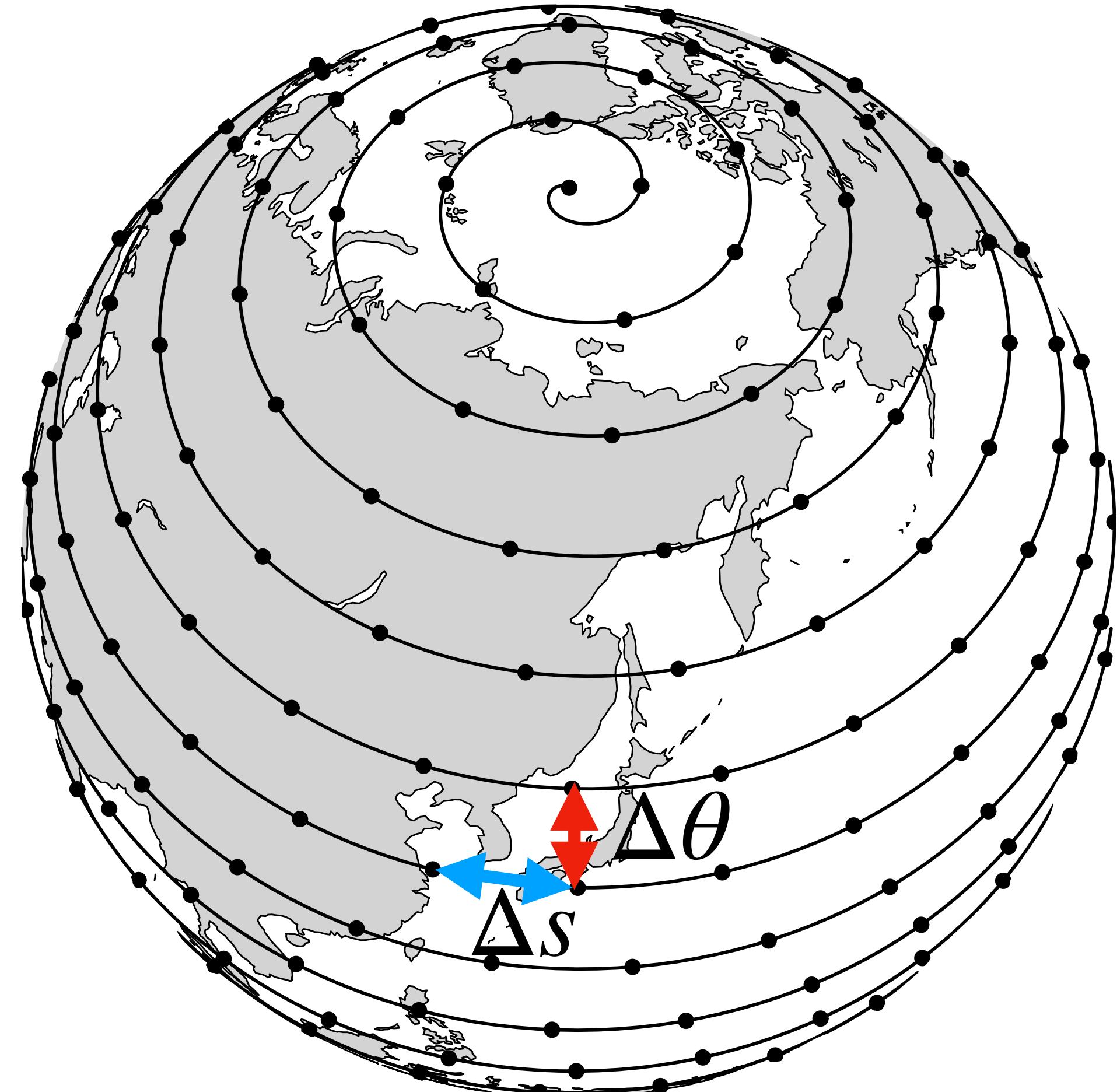
NI: Icosahedral  
NICAM



Tomita et al. 2001

# Spherical Helix

- Helix equation  $\lambda = m\theta \bmod 2\pi$
- Assume equal **interval** and **pitch** to obtain
- For almost constant intervals  $\cos \theta_k = 1 - \frac{2k-1}{n}$
- No iteration
- No limitation on the number of nodes

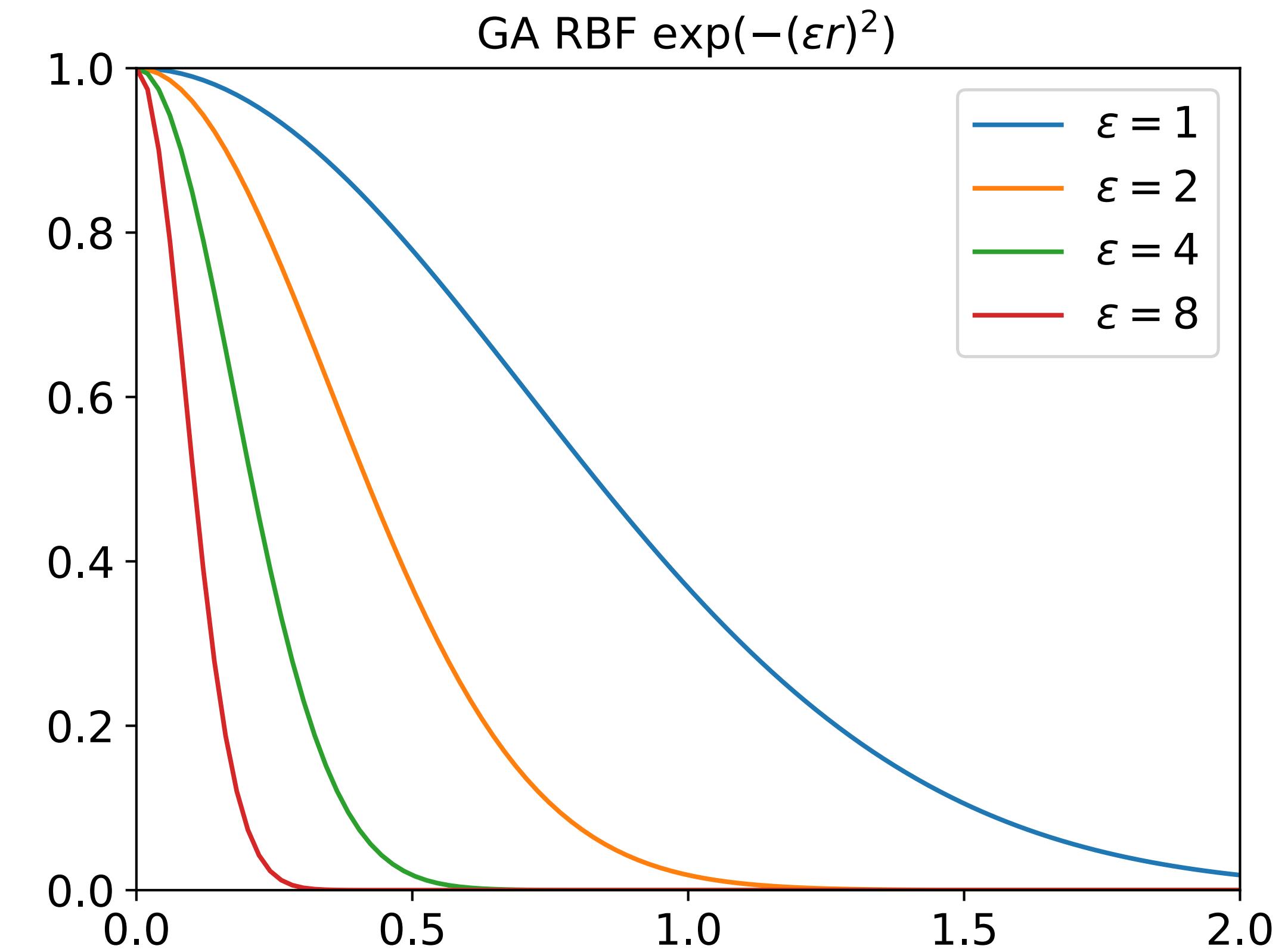


# Application of RBF to PDES

- Advection on the sphere: Flyer and Wright (2007)
- Shallow water equations: Flyer and Wright (2009)
- Mesh-free: no staggering, irregular geometry, local refinement
- Quasi-uniform nodes for a longer time step
- Spectral convergence with a very simple algorithm
- Locality: less aliasing and Gibbs phenomena (Fornberg et al. 2008)

# RBF: Radial Basis Functions

- RBF depends on the distance  $r$  from the node
- Shape parameter  $\varepsilon$  determines the radius of influence



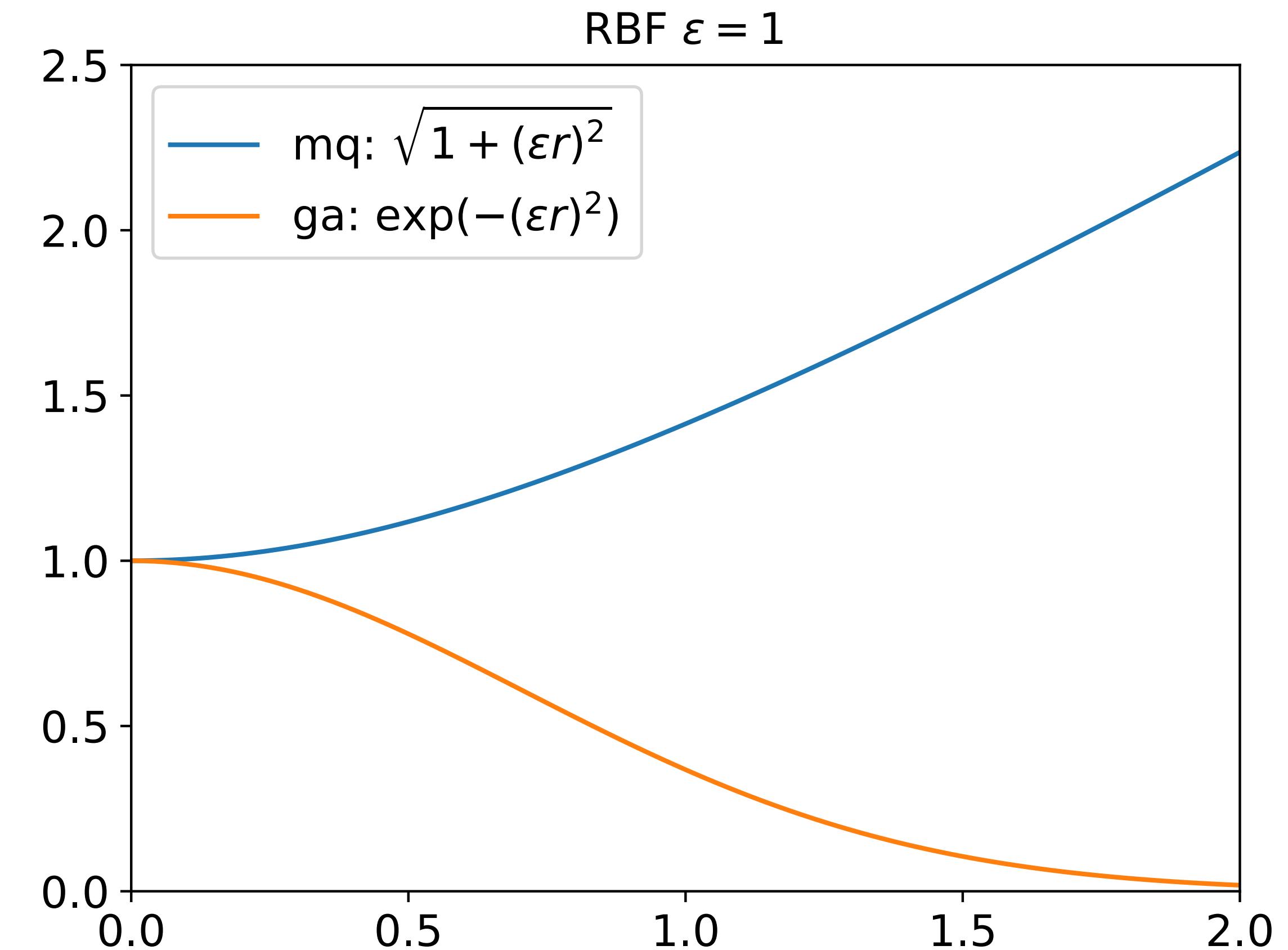
# Type of RBF

- Infinitely smooth: GA, MQ, etc

$$\frac{d}{dr} \sqrt{1 + (\varepsilon r)^2} = \frac{\varepsilon^2 r}{\sqrt{1 + (\varepsilon r)^2}}$$

$$\frac{d}{dr} \exp [-(\varepsilon r)^2] = -2\varepsilon^2 r \exp [-(\varepsilon r)^2]$$

- Piecewise smooth: cubic, thin plate spline



# RBF Interpolation

RBF expansion       $f(\underline{x}) \approx s(\underline{x}) = \sum_{k=1}^n c_k \phi(r_k) \quad r_k = \| \underline{x} - \underline{x}_k \|$

Data and nodes       $\mathbf{f} \equiv [f_1, f_2, \dots, f_n]^T \quad f_k = f(\underline{x}_k)$

Collocation condition       $\mathbf{f} = \mathbf{A}\mathbf{c} \quad \mathbf{c} \equiv [c_1, c_2, \dots, c_n]^T$

$$\mathbf{A} \equiv \begin{bmatrix} \phi(0) & \phi(r_{1,2}) & \cdots & \phi(r_{1,n}) \\ \phi(r_{2,1}) & \phi(0) & \cdots & \phi(r_{2,n}) \\ \vdots & \vdots & \ddots & \vdots \\ \phi(r_{n,1}) & \phi(r_{n,2}) & \cdots & \phi(0) \end{bmatrix} \quad r_{i,j} \equiv = \| \underline{x}_i - \underline{x}_j \|$$

# Nodal weight

- Surface of a unit sphere
- uniform weight of a node
- Sums of each column of RBF interpolation matrix

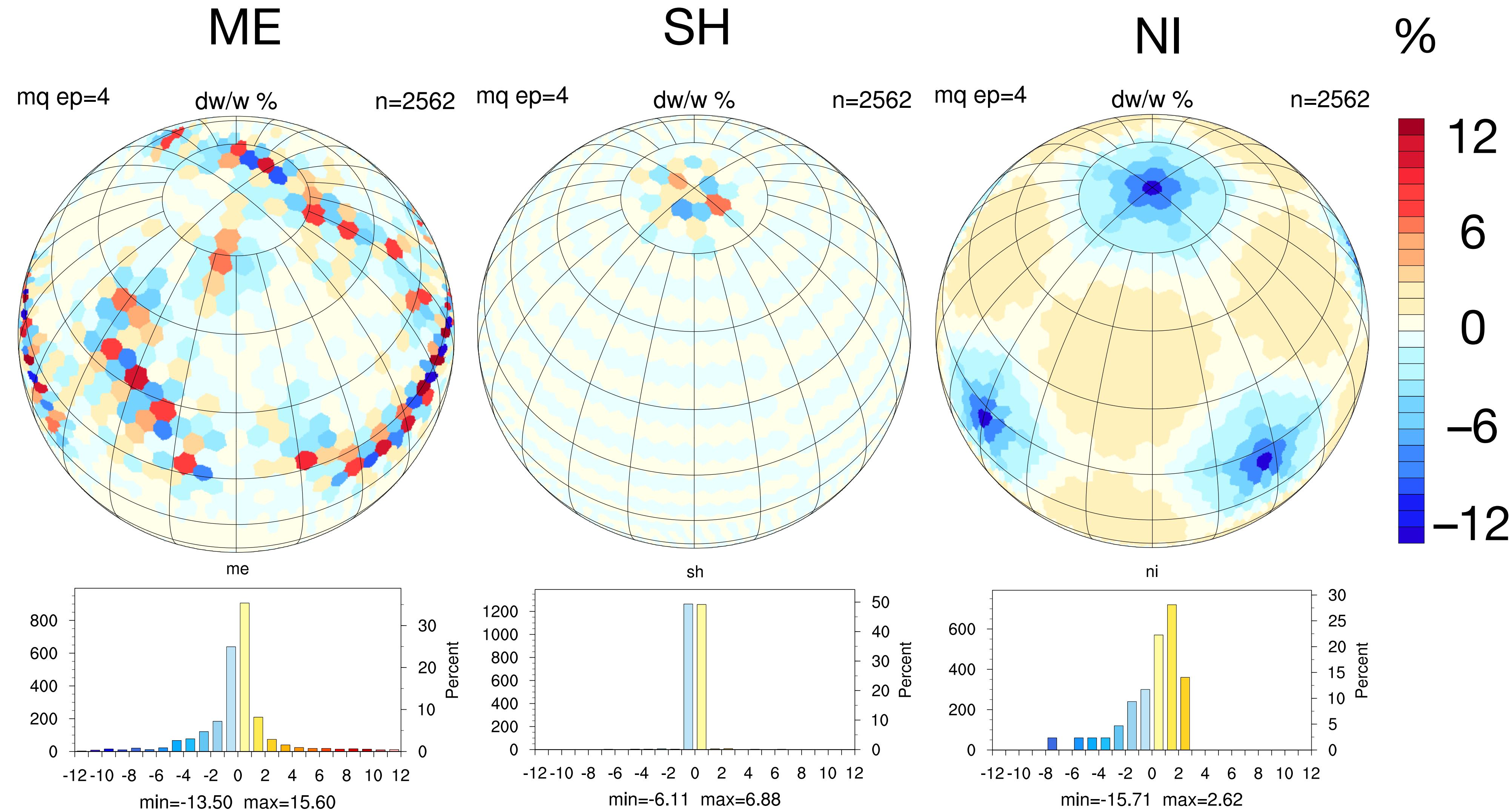
$$4\pi$$

$$4\pi/n$$

$$\mathbf{w} = 4\pi \frac{\mathbf{A}^{-1}\mathbf{e}}{\mathbf{e}^T \mathbf{A}^{-1} \mathbf{e}}$$

$$\mathbf{e} \equiv [1, 1, \dots, 1]^T$$

# Error from uniform weight $4\pi/n$



# Derivative operator

Distance       $r_k(\underline{x}) = \|\underline{x} - \underline{x}_k\| = \sqrt{(x - x_k)^2 + (y - y_k)^2 + (z - z_k)^2}$

Differentiate RBF       $\nabla f(\underline{x}) \approx \nabla \left[ \sum_{k=1}^n c_k \phi(r_k(\underline{x})) \right] = \sum_{k=1}^n c_k \nabla \phi(r_k(\underline{x}))$

Gradient of RBF       $\nabla \phi(r_k(\underline{x})) = \frac{d\phi(r_k(\underline{x}))}{dr} \nabla r_k(\underline{x}) = \phi'(r_k(\underline{x})) \frac{\underline{x} - \underline{x}_k}{r_k(\underline{x})}$

# Shallow water model

Cartesian formulation

$$\frac{\partial \underline{u}}{\partial t} = - (\underline{u} \cdot P \nabla) \underline{u} - f(\underline{x} \times \underline{u}) - g \nabla h + \mu \underline{x}$$

$$\frac{\partial h}{\partial t} = - \nabla \cdot (h \underline{u})$$

$P \equiv I - \underline{x} \underline{x}^T$  projects a vector to the tangent plane

$\mu = \underline{x} \cdot \left\{ (\underline{u} \cdot P \nabla) \underline{u} \right\}$  Lagrange multiplier (Côté 1988)

# Projection to the tangent plane

radial component  $\underline{x}\underline{x}^\top \underline{u}$

tangential component  $\underline{u} - \underline{x}\underline{x}^\top \underline{u}$

$$P = I - \underline{x}\underline{x}^\top = \begin{bmatrix} 1 - x^2 & -xy & -xz \\ -xy & 1 - y^2 & -yz \\ -xz & -yz & 1 - z^2 \end{bmatrix}$$

$$P \nabla \phi \left( r_k \left( \underline{x} \right) \right) = \begin{bmatrix} \underline{x}\underline{x}^\top \underline{x}_k - \underline{x}_k \\ \underline{y}\underline{x}^\top \underline{x}_k - \underline{y}_k \\ \underline{z}\underline{x}^\top \underline{x}_k - \underline{z}_k \end{bmatrix} \frac{\phi' \left( r_k \left( \underline{x} \right) \right)}{r_k \left( \underline{x} \right)}$$

# Derivative matrix

$$P \nabla \mathbf{f} \Big|_x = (\mathbf{B}^x \mathbf{A}^{-1}) \mathbf{f}$$

$$\mathbf{B}_{j,k}^k \equiv \left[ x_j x_{-j-k}^\top - x_k \right] \frac{\phi'(r_{j,k})}{r_{j,k}}$$

Solve the linear system for  $\mathbf{D}^x$

$$\mathbf{D}^x \mathbf{A} = \mathbf{B}^x$$

Similarly,  $\mathbf{D}^y$  and  $\mathbf{D}^z$  are obtained.

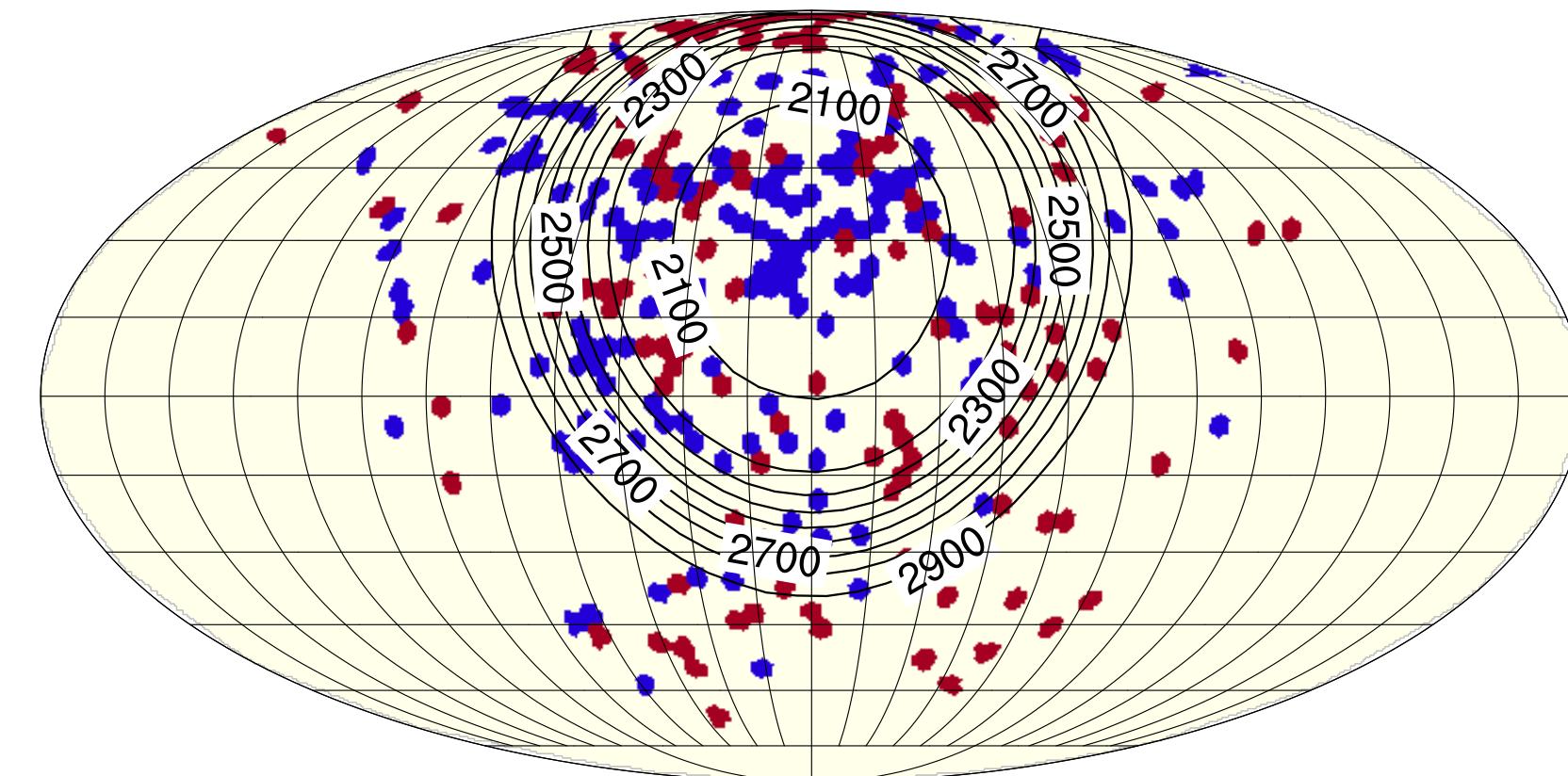
# Numerical studies

- Standard test for shallow water equations (Williamson et al. 1992)
  - Seven tests including steady-state, flow over isolated mountain, and Rossby-Haurwitz wave
  - RBF filter (Fornberg and Lehto 2011) as numerical diffusion in Case 5–7
- 4th Runge-Kutta for time integration
- Compare ME, NI, and SH with 2562 nodes (g-level 4)~T51

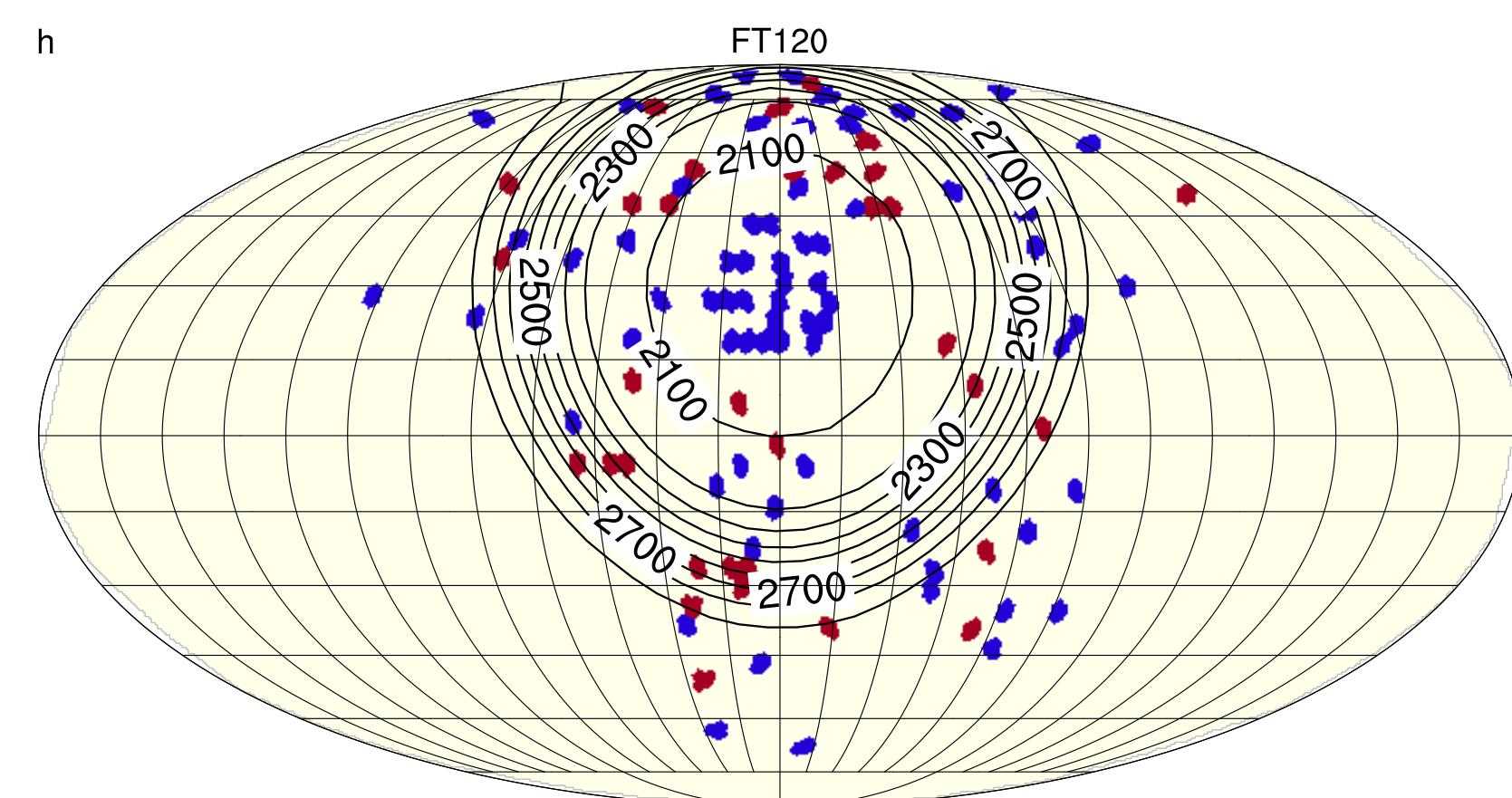
# Case 3: Geostrophic flow with compact support

$n = 2562; \Delta t = 24 \text{ min } \varepsilon = 4.75$

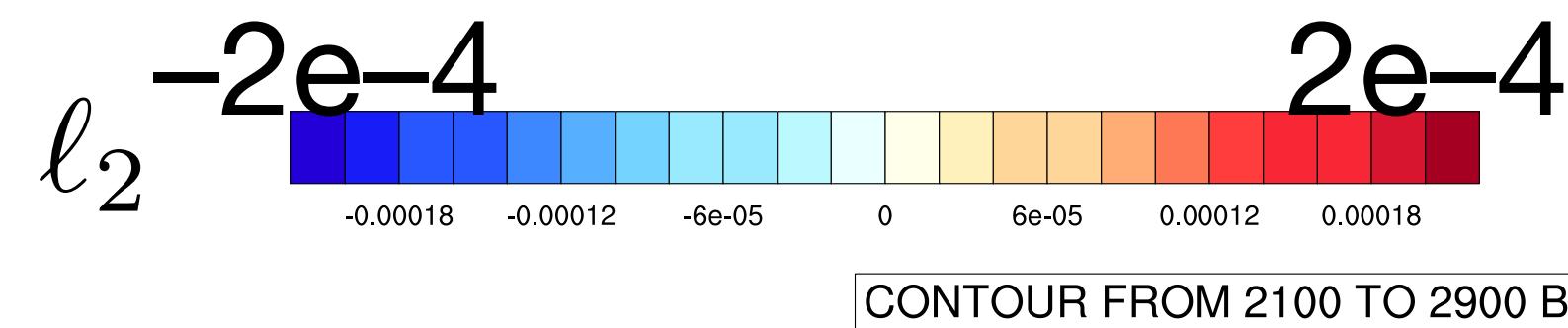
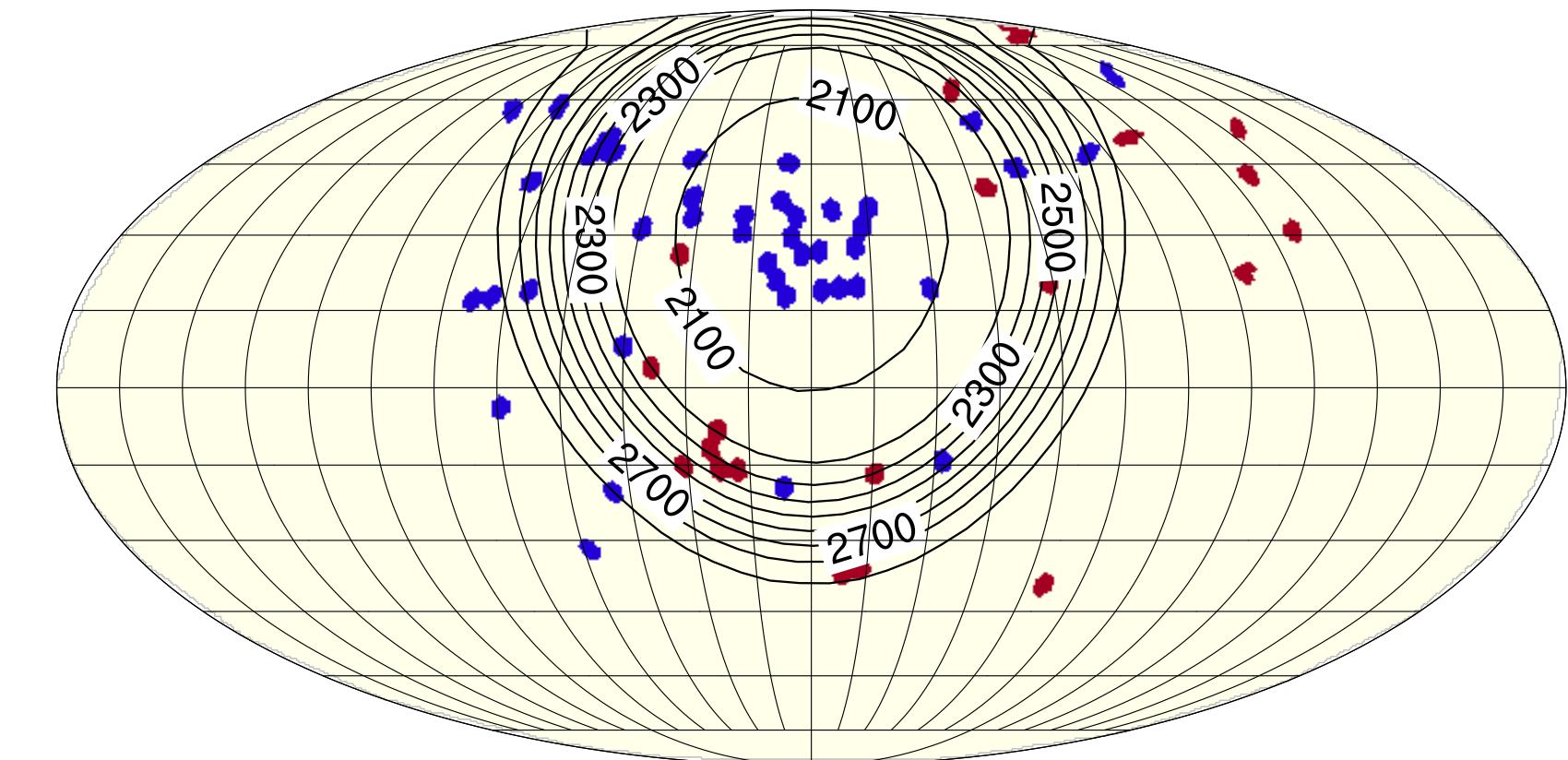
ME



SH

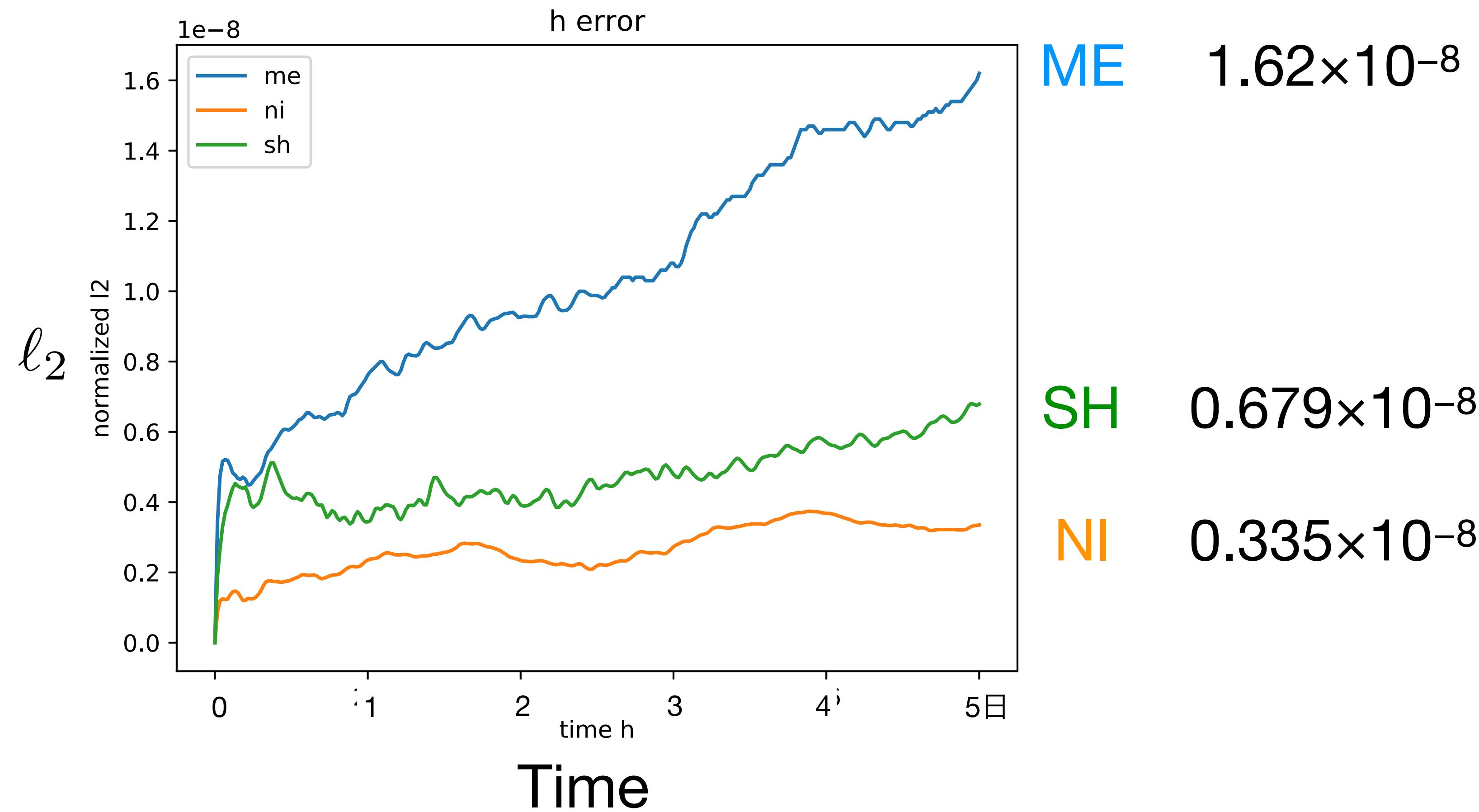


NI



# Case 3: Geostrophic flow with compact support

$n = 2562; \Delta t = 24 \text{ min } \varepsilon = 4.75$



method		no. of nodes ( $N$ )	time step ( $\Delta t$ )	relative $\ell_2$ error in $h$
RBF	<b>leapfrog</b>	784 (28)	20 min	$6.32 \times 10^{-6}$
		1849 (43)	12 min	$1.97 \times 10^{-8}$
<b>Flyer and Wright 2009</b>		3136 (56)	10 min	$3.65 \times 10^{-10}$
		4096 (64)	8 min	$4.72 \times 10^{-11}$
<b>Spherical Harmonics</b> SH; Jakob-Chien <i>et al.</i> (1995)		5041 (71)	6 min	$6.88 \times 10^{-12}$
		8192 (1849)	20 min* (3)	$7 \times 10^{-10}$
<b>Double Fourier</b> DF/SHF; Spotz <i>et al.</i> (1998)		18 432 (4096)	15 min*	$2.5 \times 10^{-10}$
		2048	6 min	$2 \times 10^{-6}$
<b>Spectral Element</b> SE; Taylor <i>et al.</i> (1997)		8192	3 min	$4 \times 10^{-10}$
		32 768	90 s	$2 \times 10^{-13}$
		6144	90 s	$8 \times 10^{-7}$
		24 576	45 s	$1 \times 10^{-10}$
RBF	SH	RK4	2562 (51)	24 min
				$6.79 \times 10^{-9}$

\*Semi-implicit

# RBF hyperviscosity

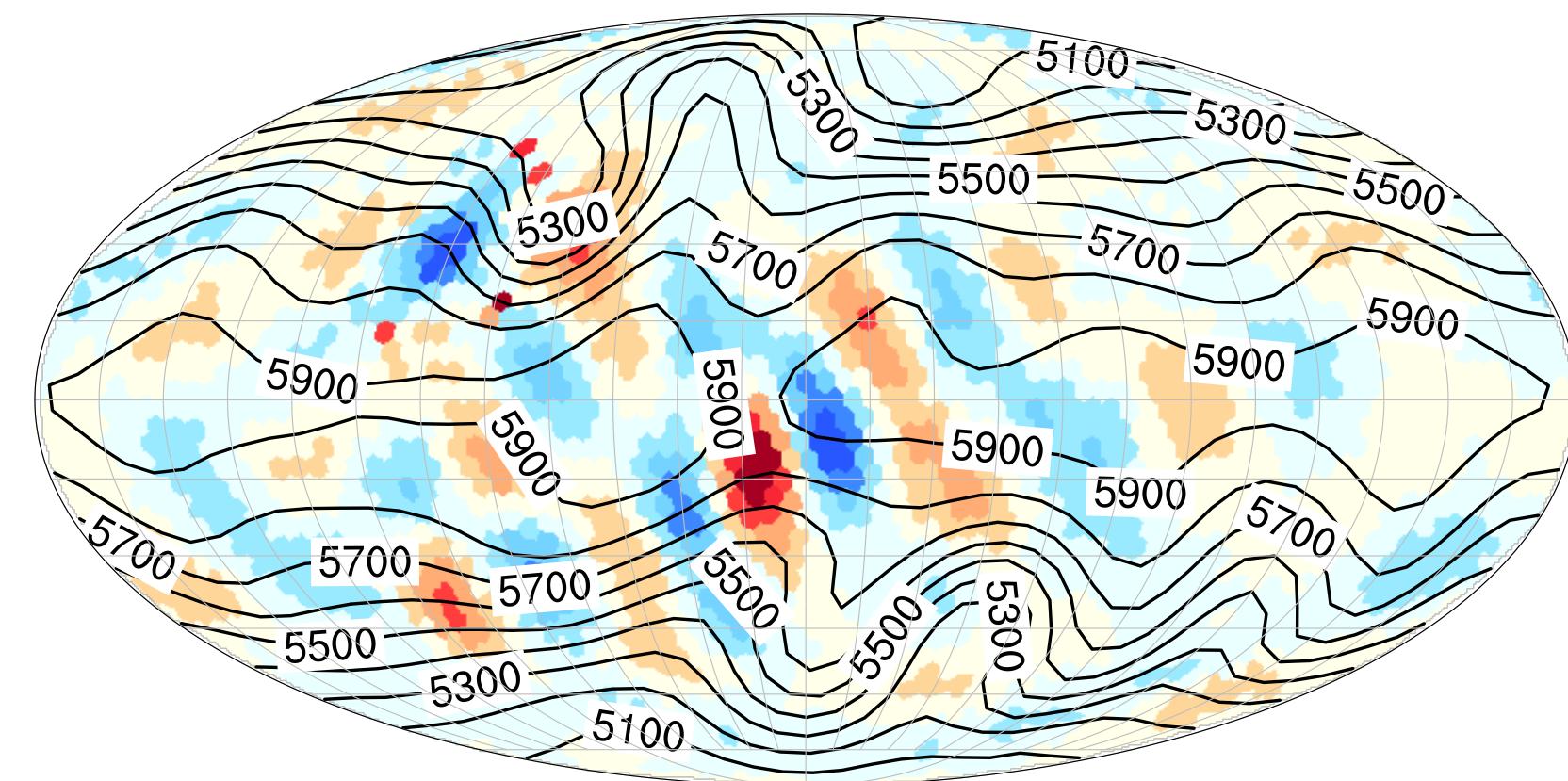
Fornberg and Lehto 2011

- $A^{-1}$  damps low order eigenfunctions weakly and high order eigenfunctions extremely fast.
- Irrespective of node distribution, number of dimensions, type of RBF.
- No decision required for what power of the Laplacian.
- No ode can grow for any  $\epsilon$

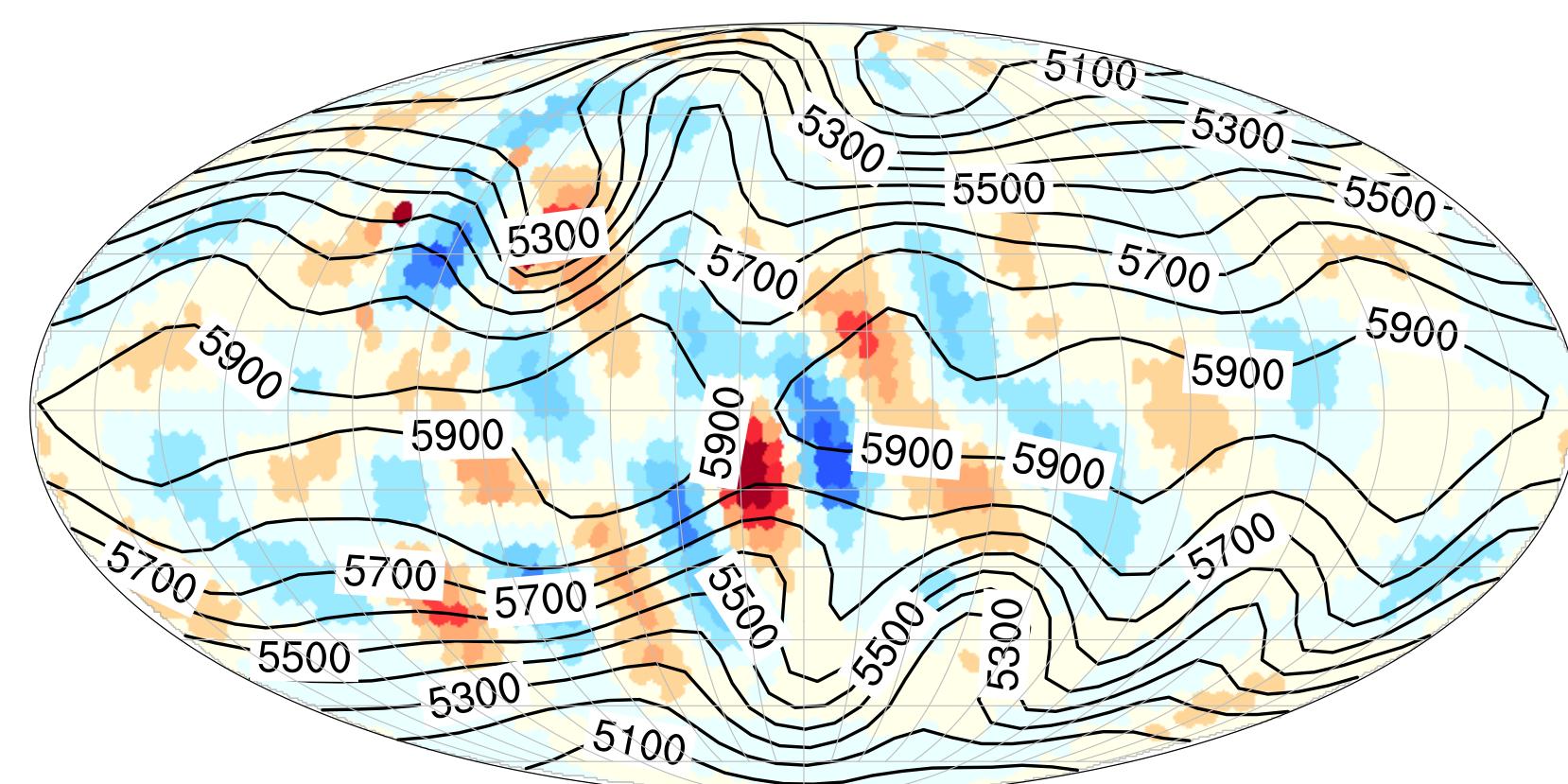
# Case 5: Flow over an isolated mountain

$$n = 2562 \Delta t = 15 \text{ min } \varepsilon = 4.25 \gamma = -2 \times 10^{-8}$$

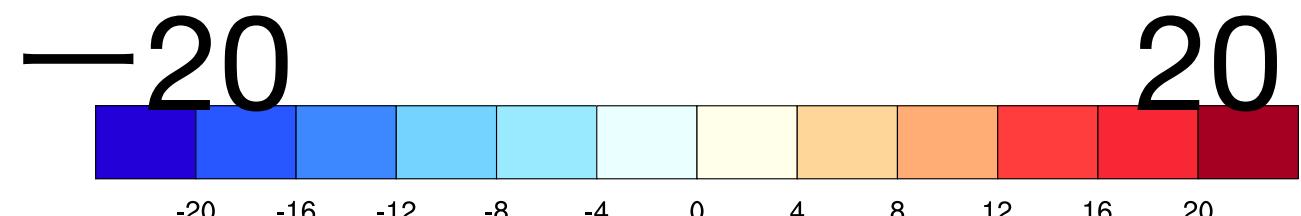
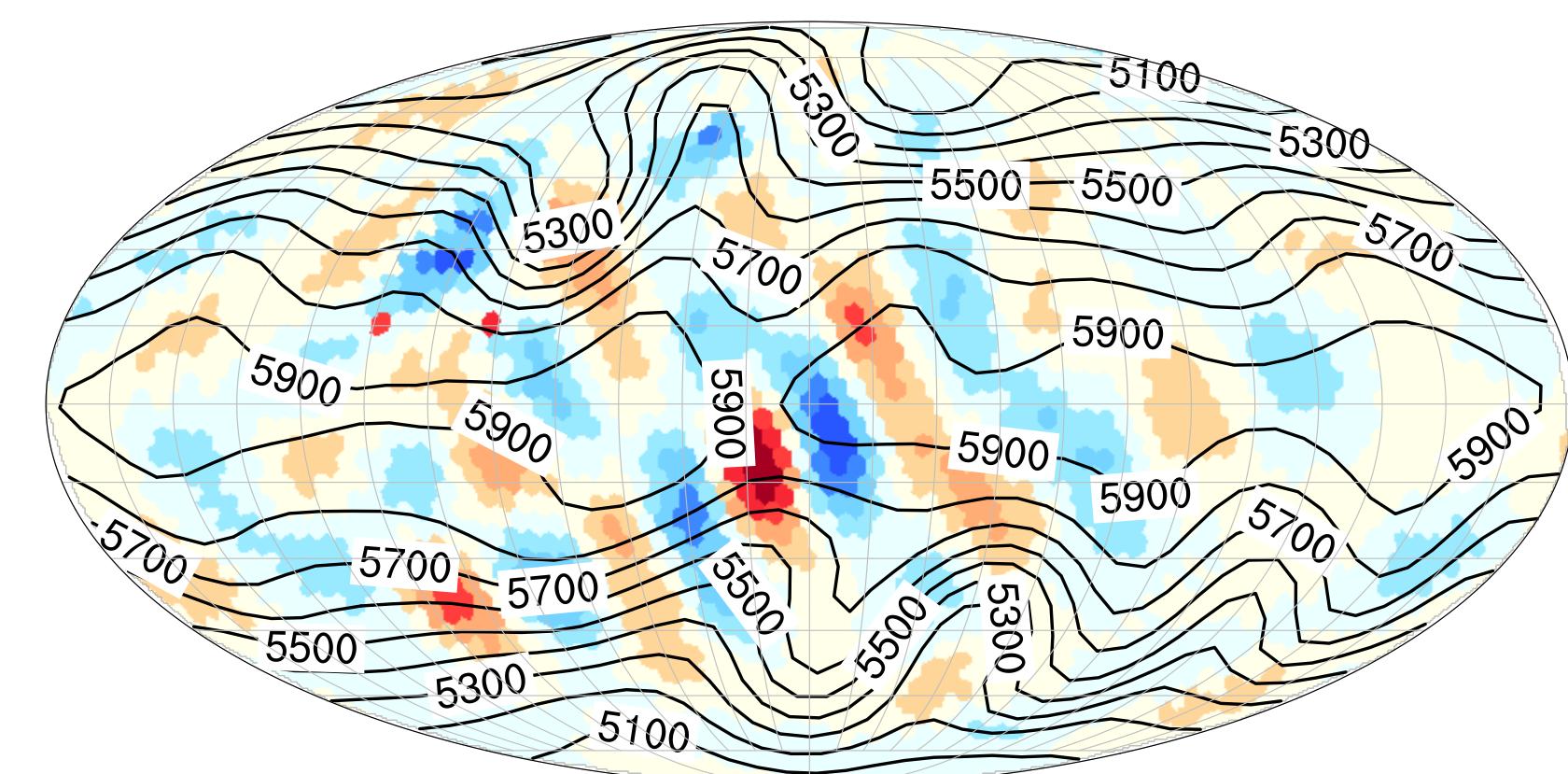
ME



SH



NI

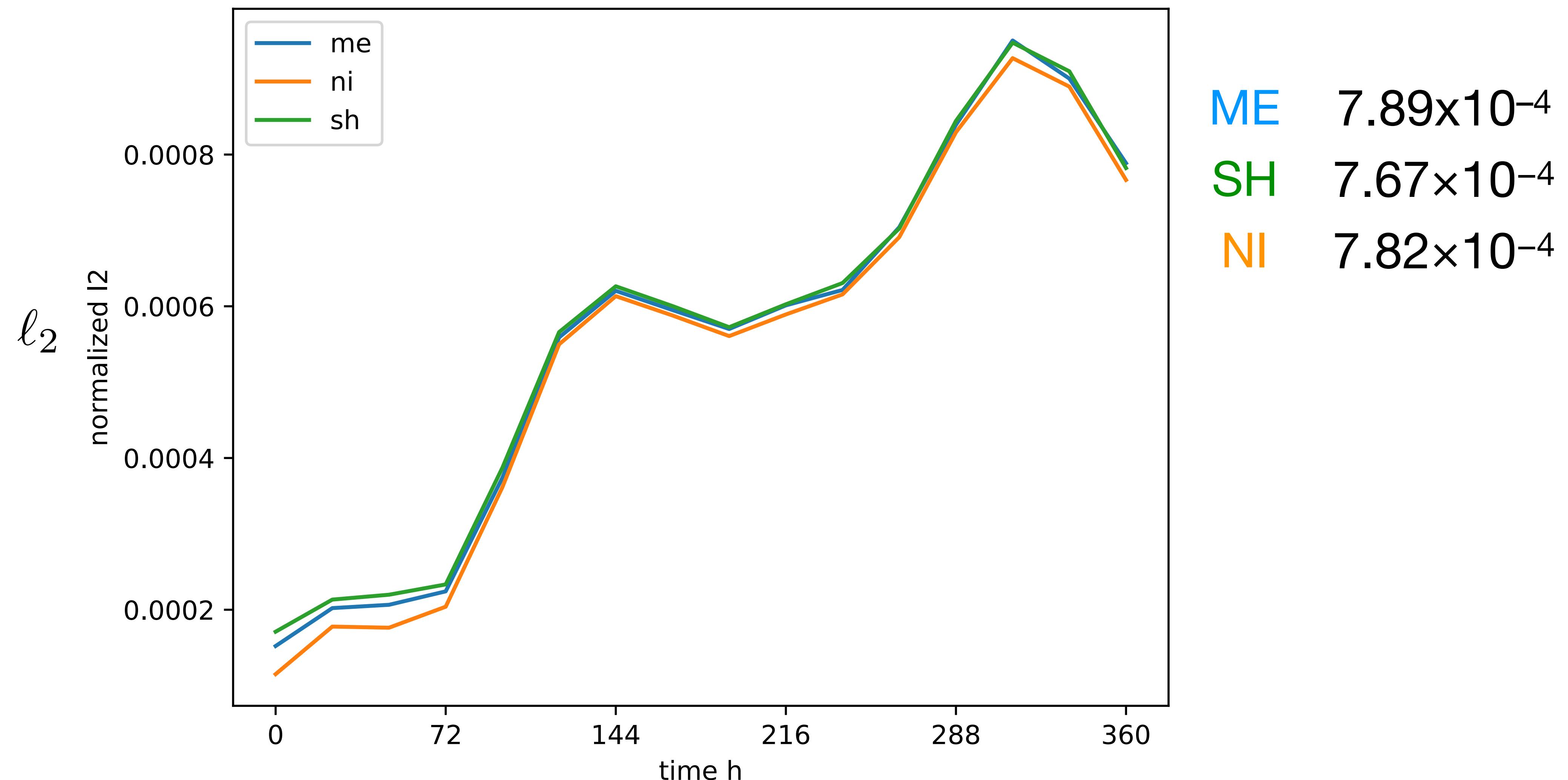


CONTOUR FROM 5100 TO 5900 BY 100

# Case 5: Flow over an isolated mountain

$$n = 2562 \Delta t = 15 \text{ min } \varepsilon = 4.25 \gamma = -2 \times 10^{-8}$$

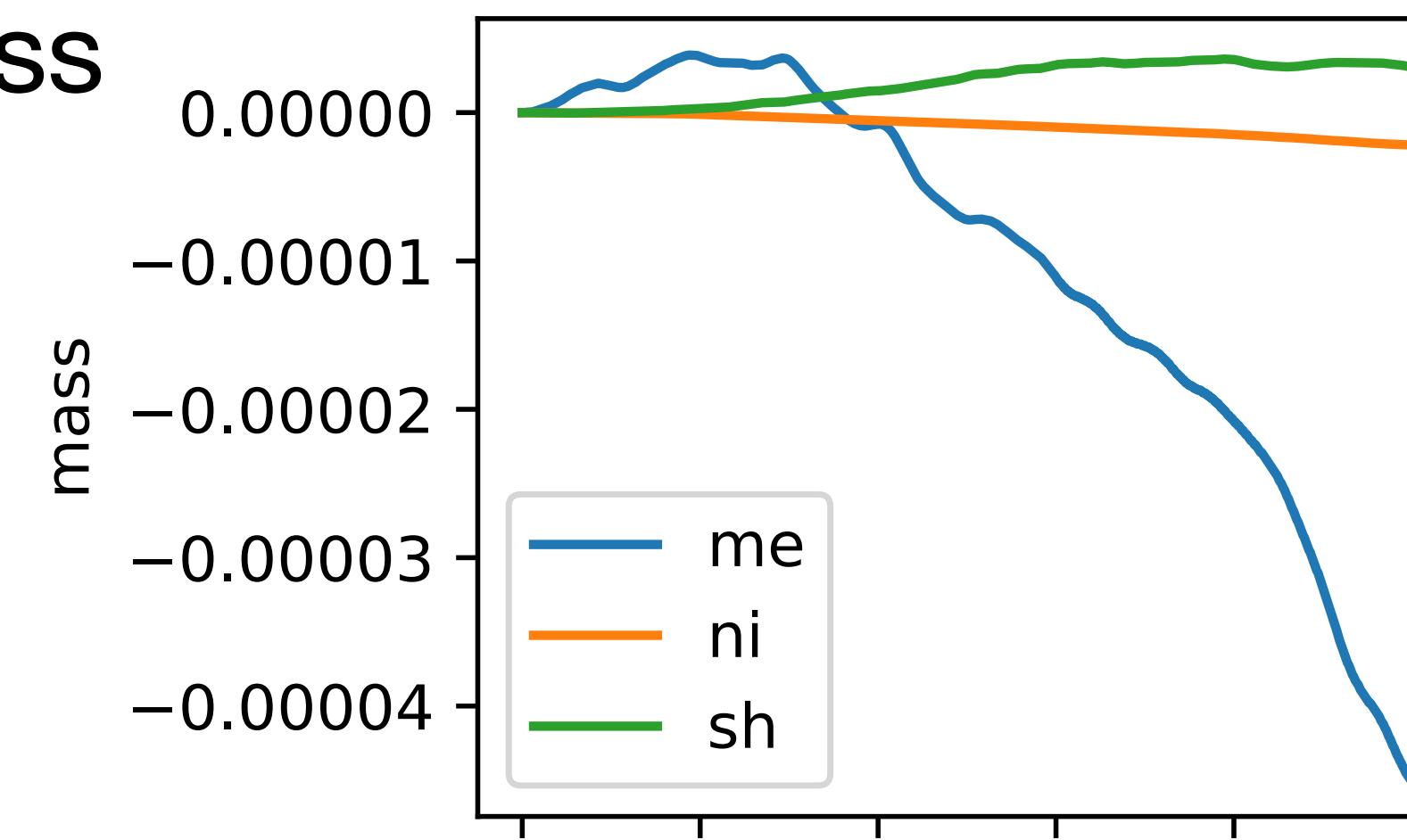
h error from T213



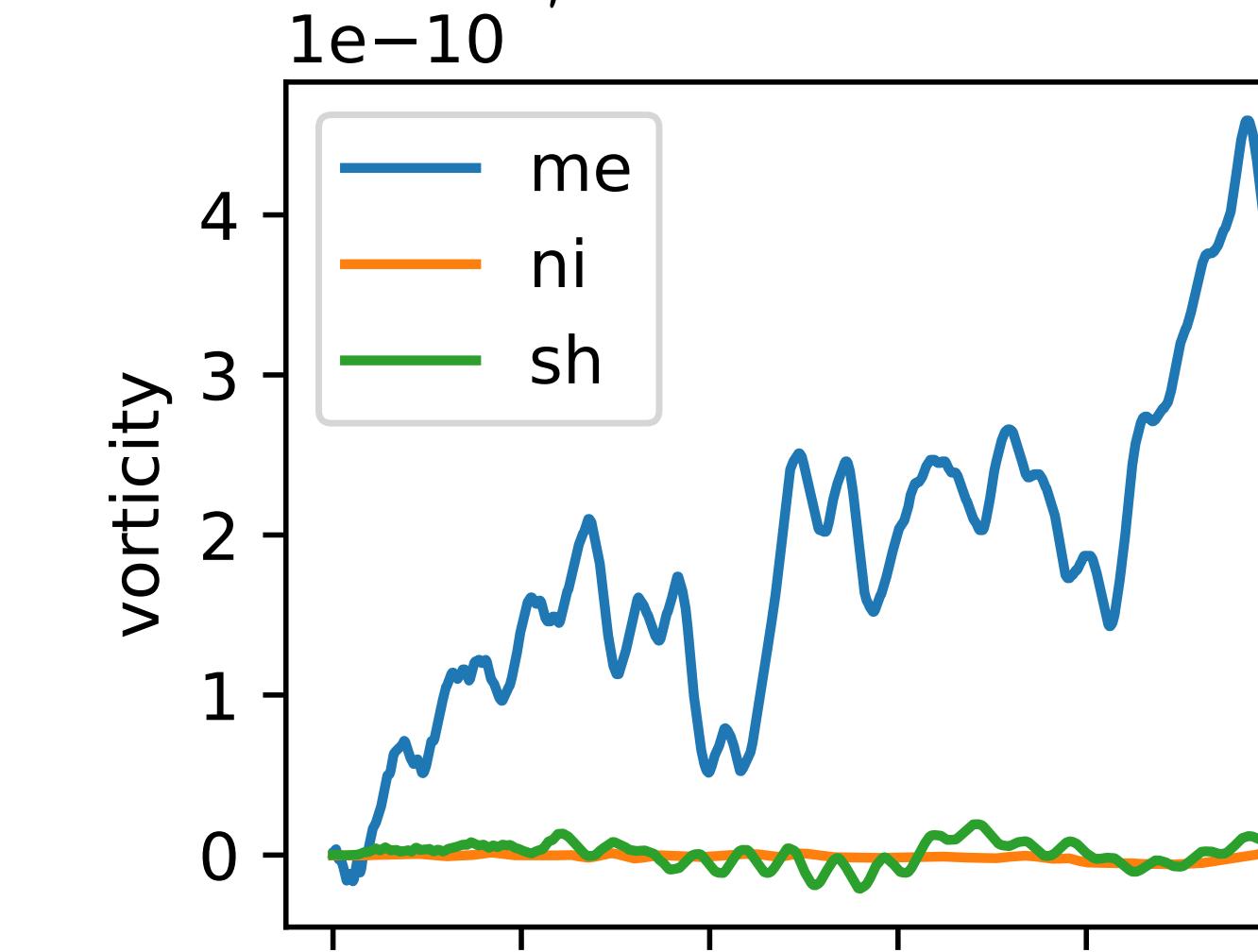
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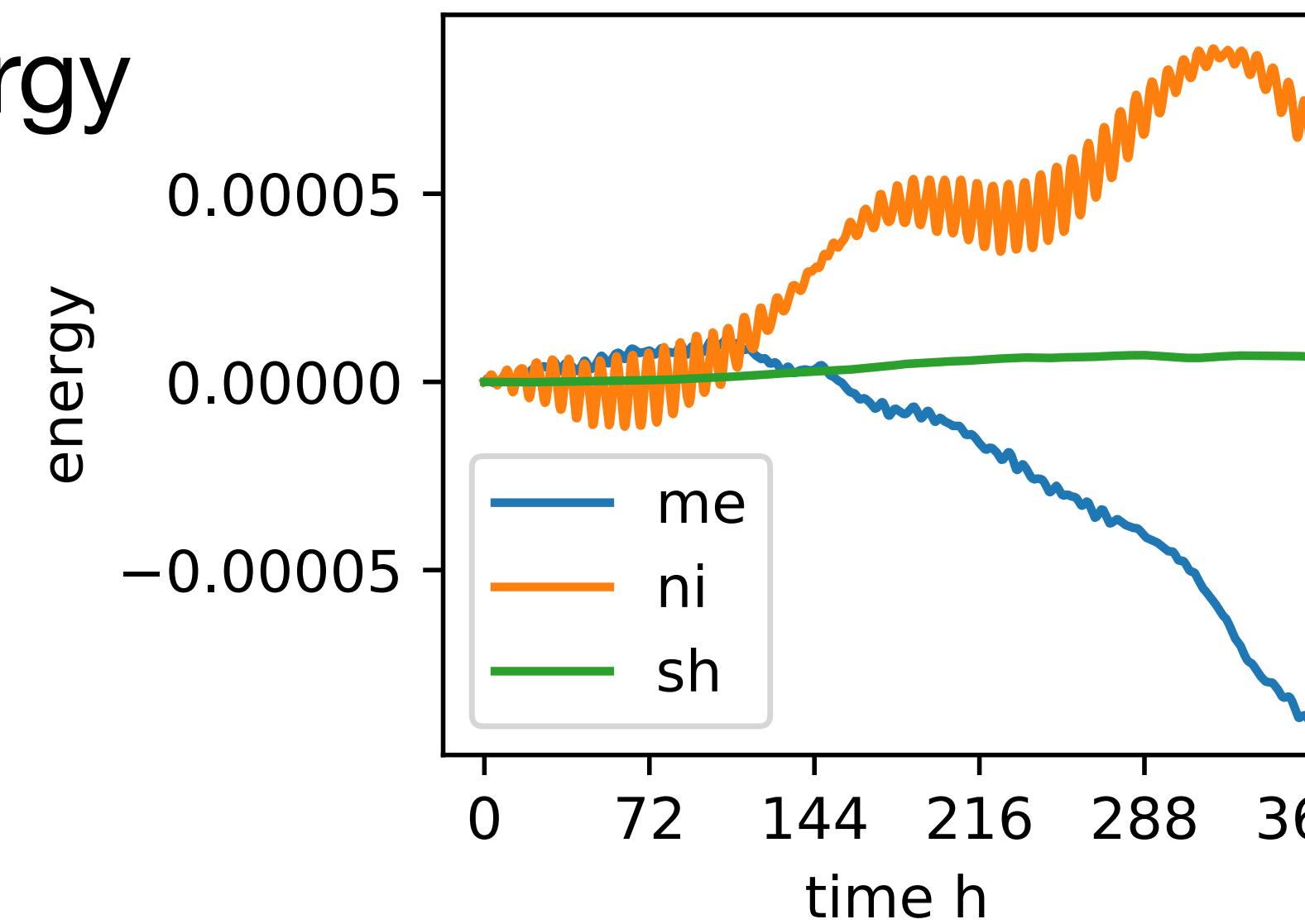
mass



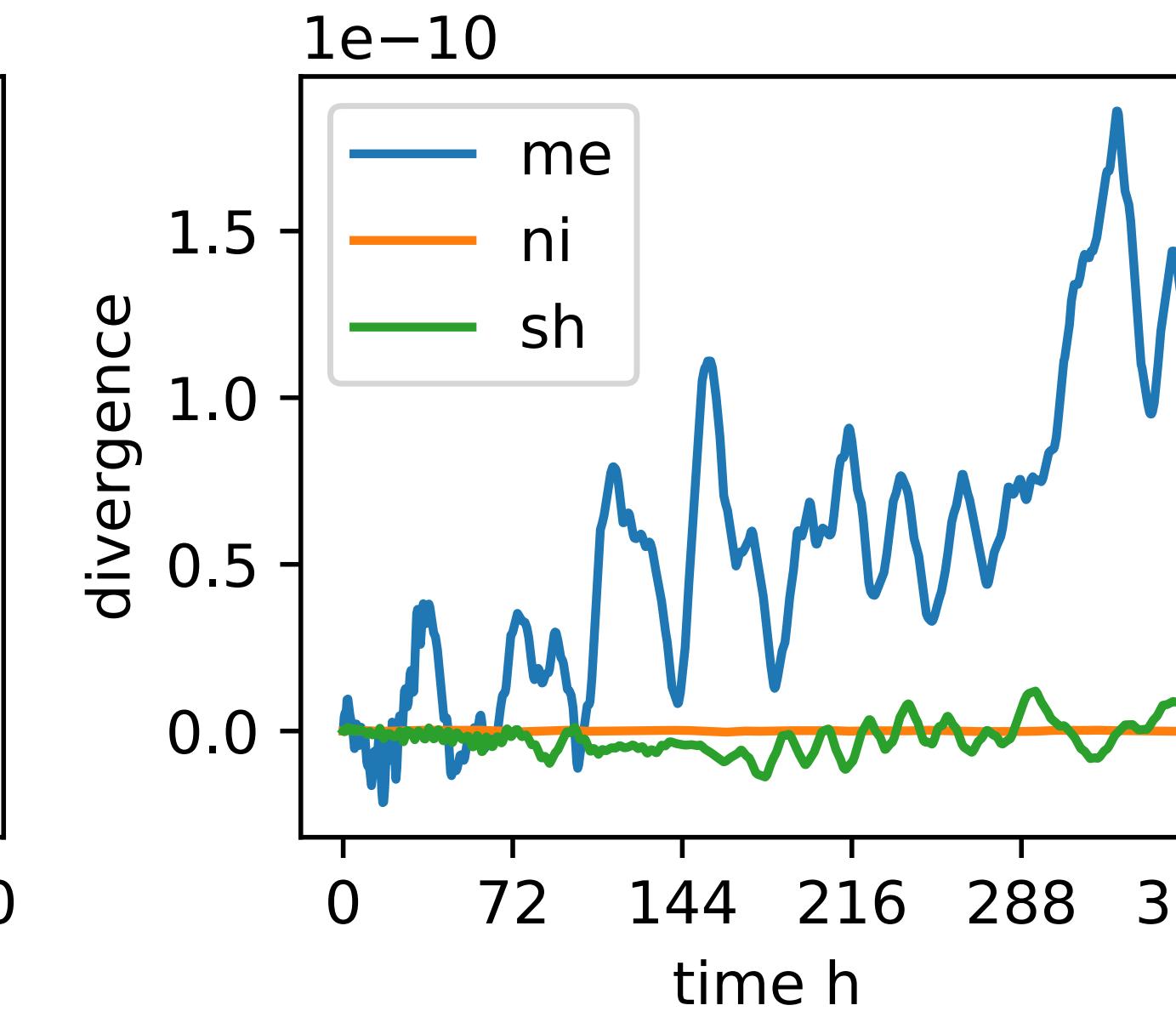
vorticity



energy



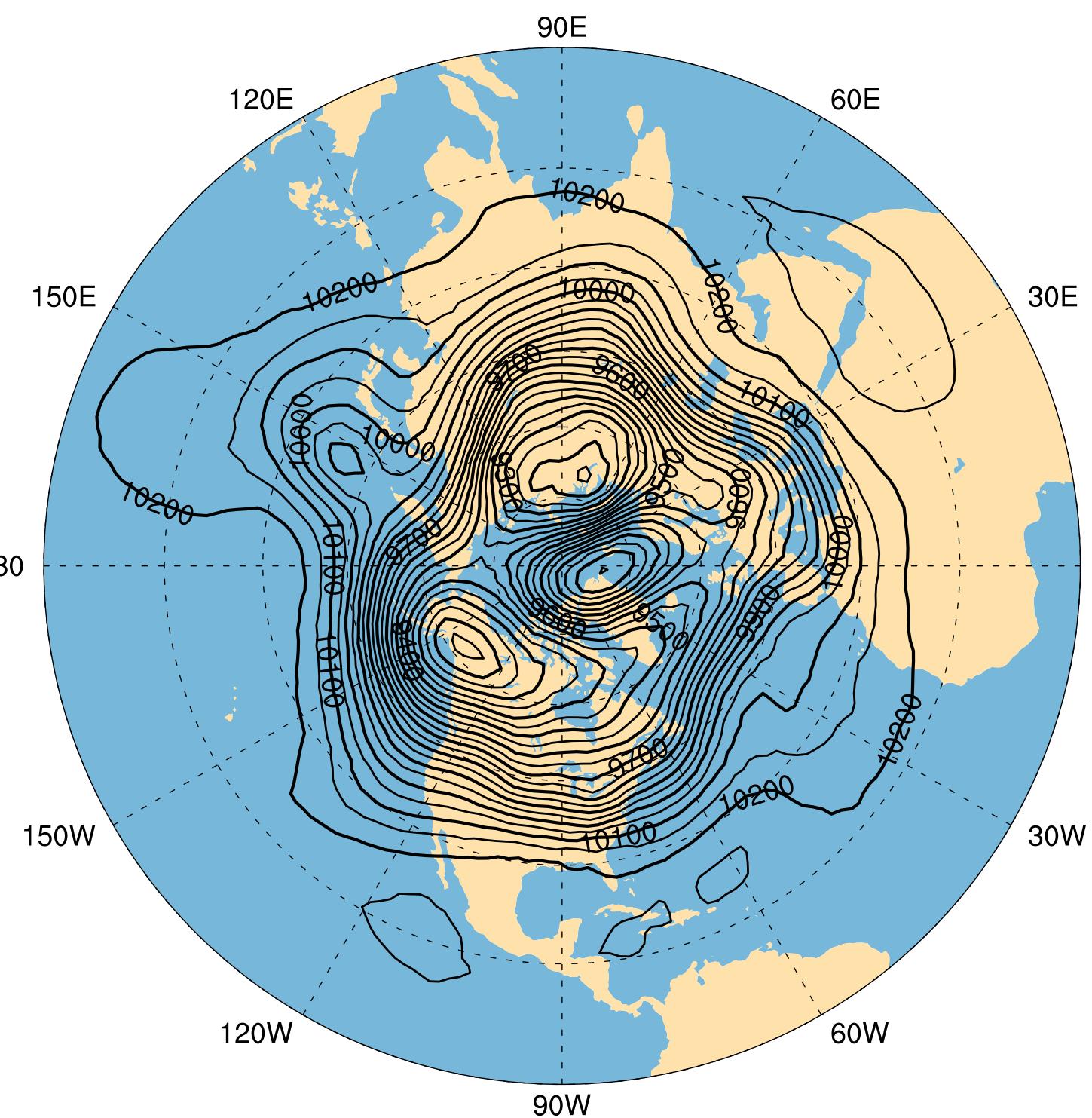
divergence



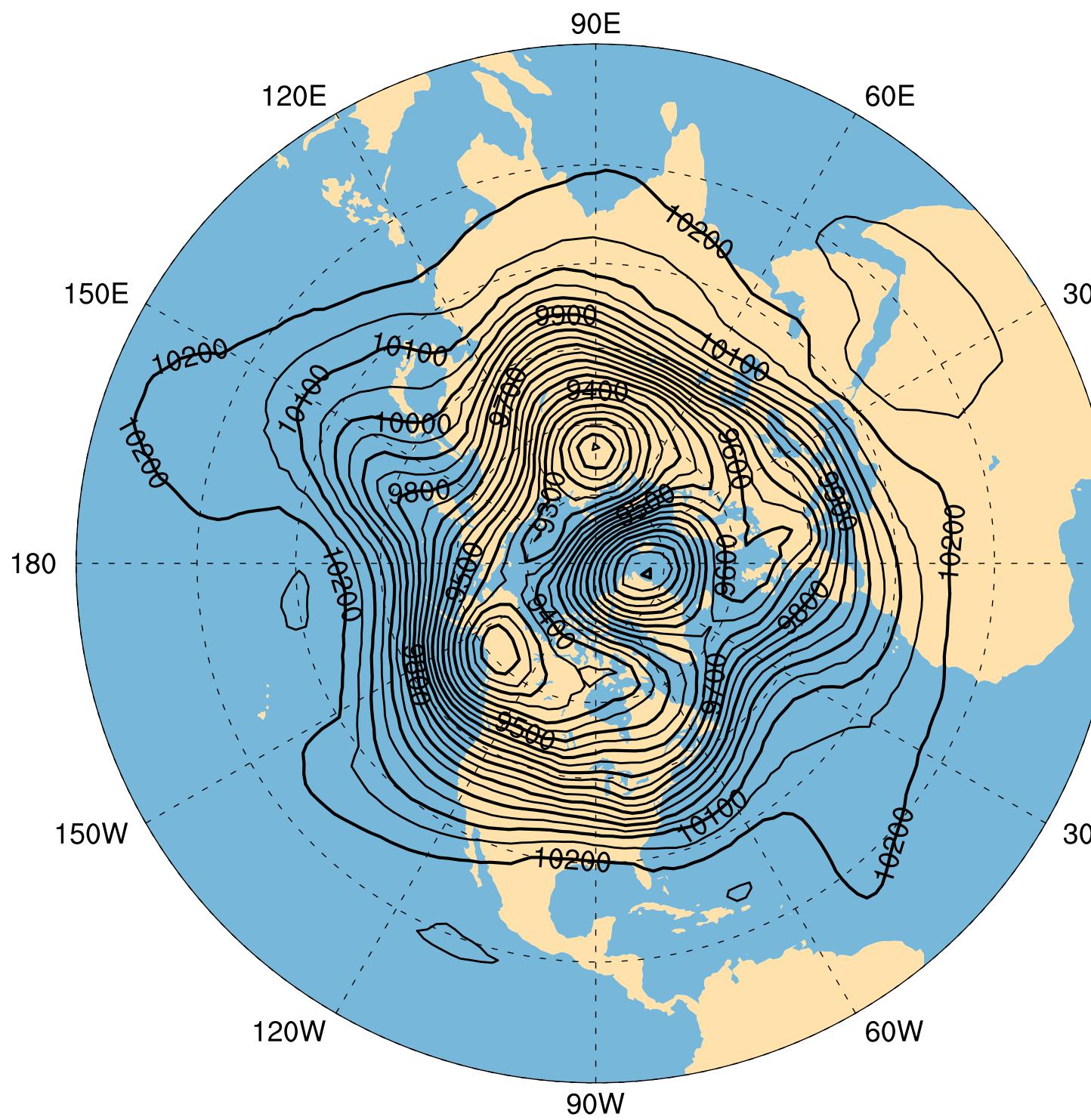
# Case 7: Forecast from 500 hPa analysis

$$n = 2562 \Delta t = 15 \text{ min } \varepsilon = 4.25 \gamma = -2 \times 10^{-7} \quad t = 5 \text{ 日}$$

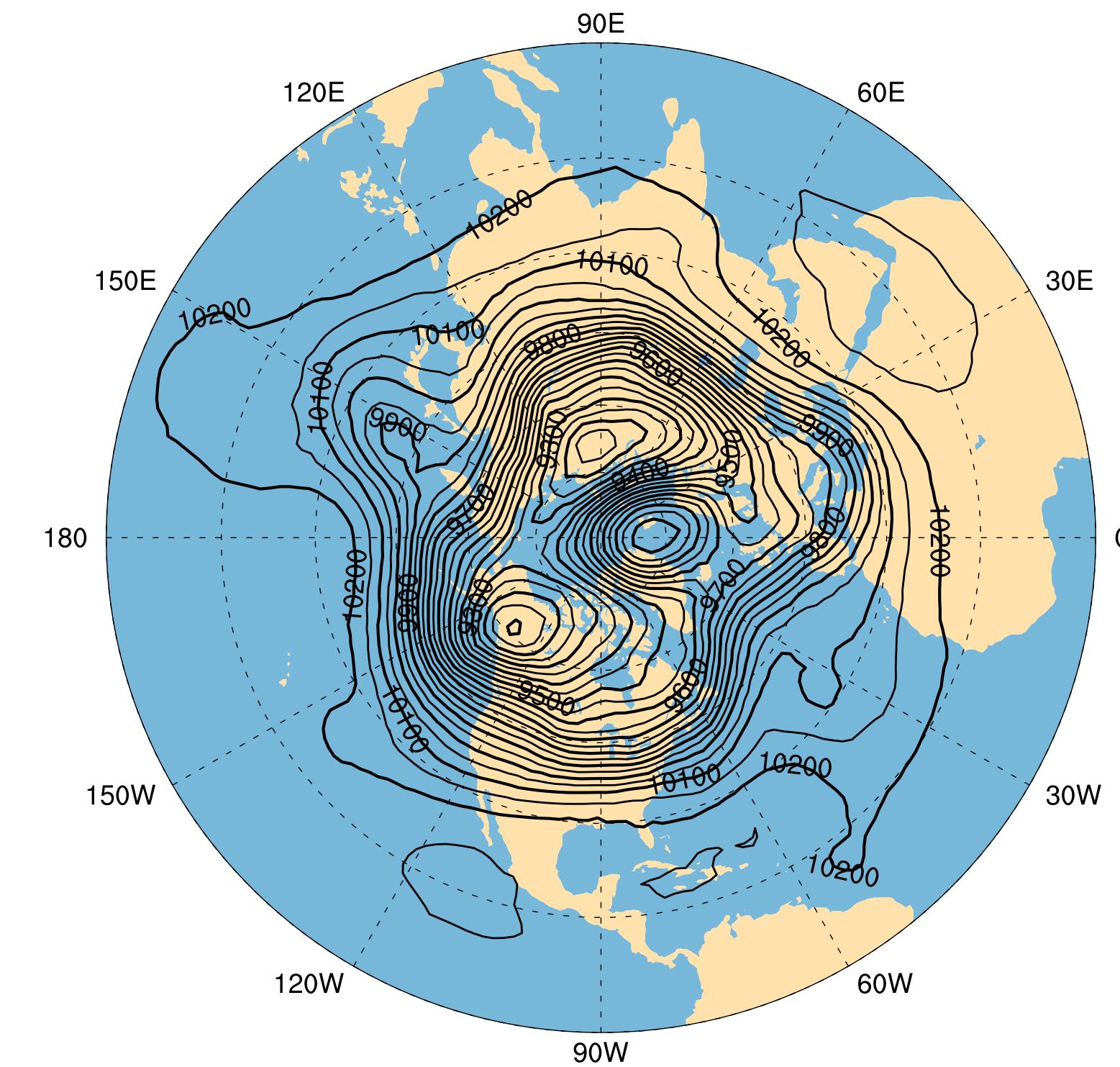
ME



SH



NI



CONTOUR FROM 9050 TO 10250 BY 50

CONTOUR FROM 9050 TO 10250 BY 50

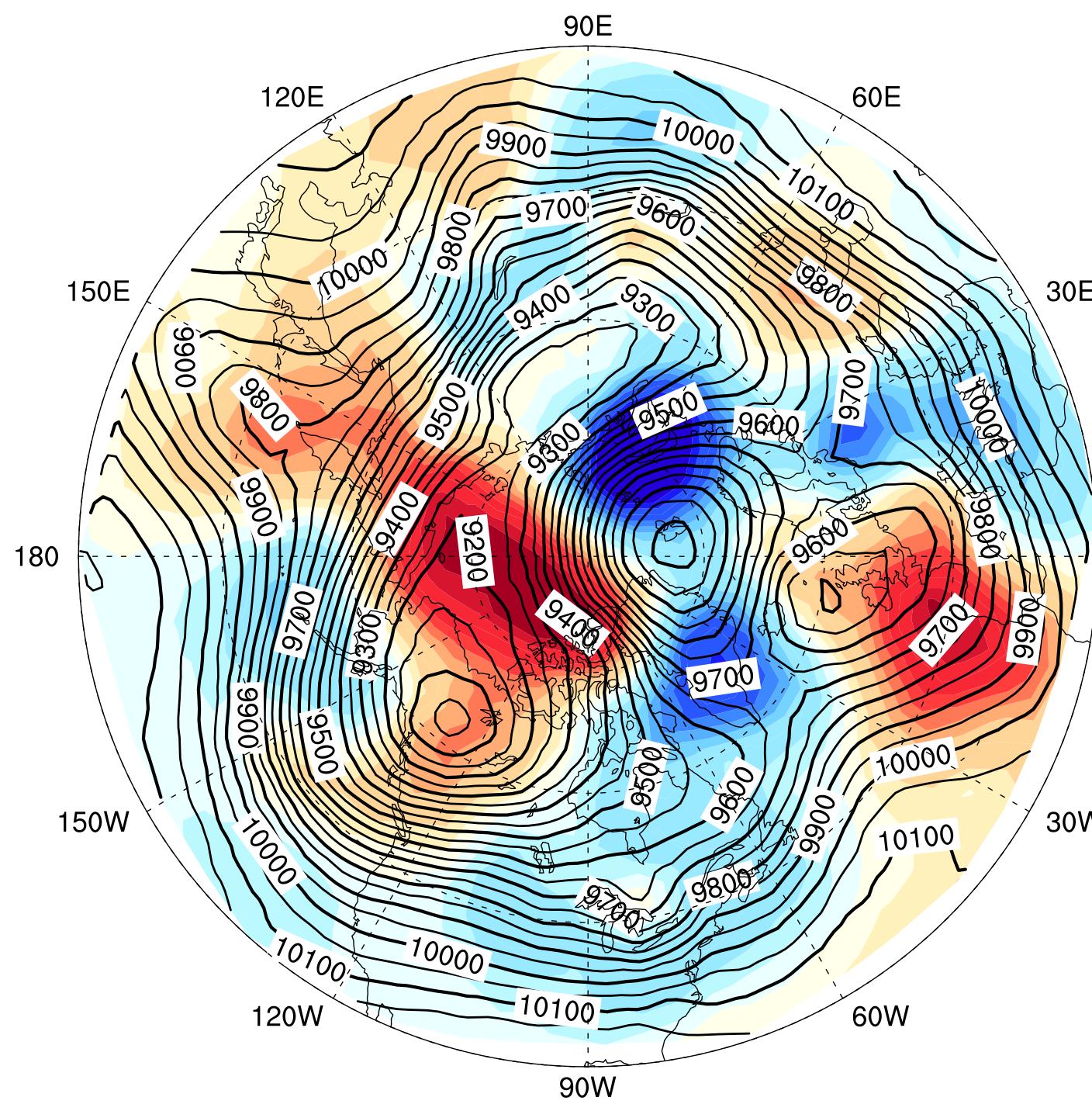
CONTOUR FROM 9050 TO 10250 BY 50

# Case 7: Forecast from 500 hPa analysis

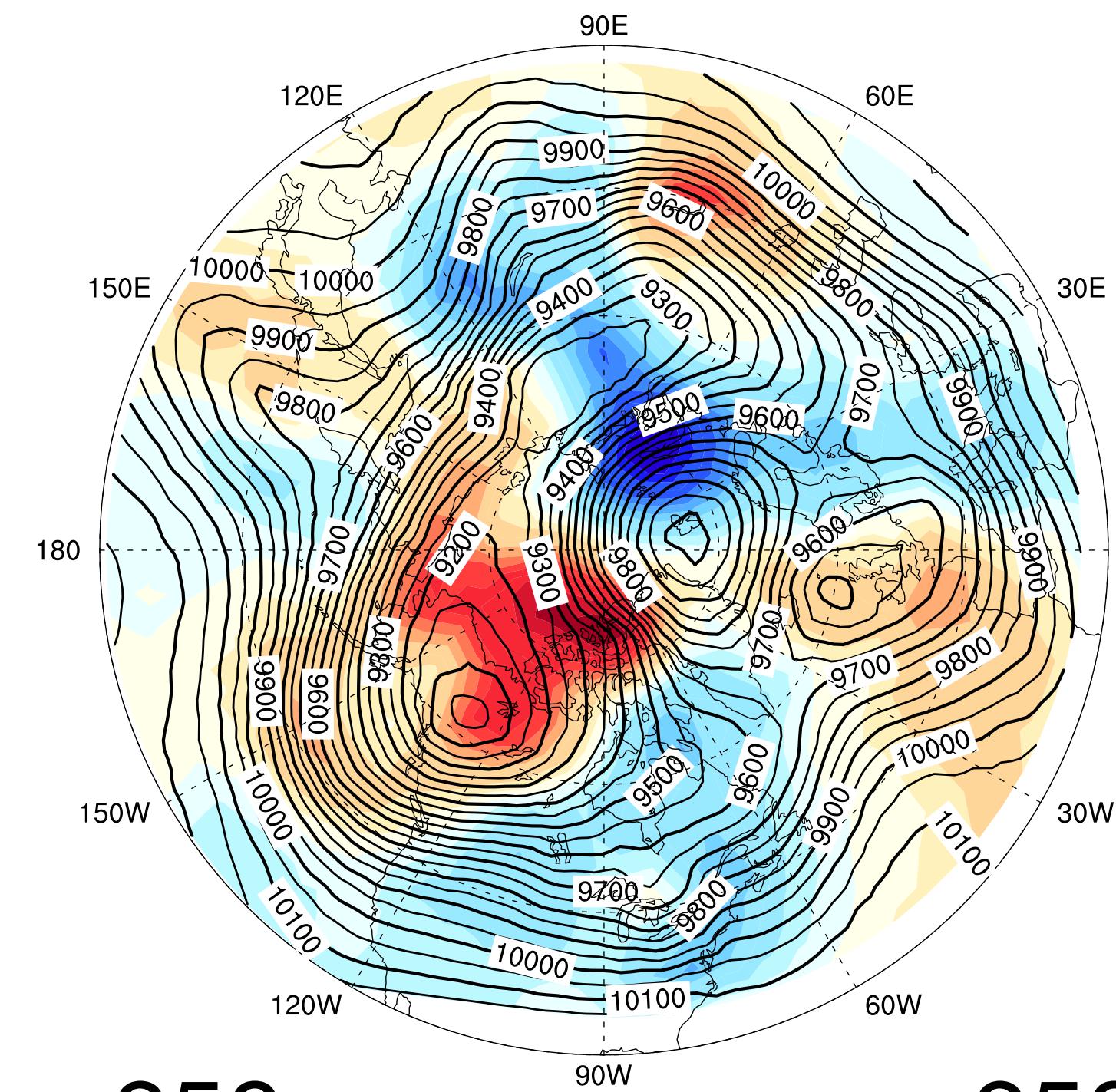
Error from reference T213 (contours)

>30N

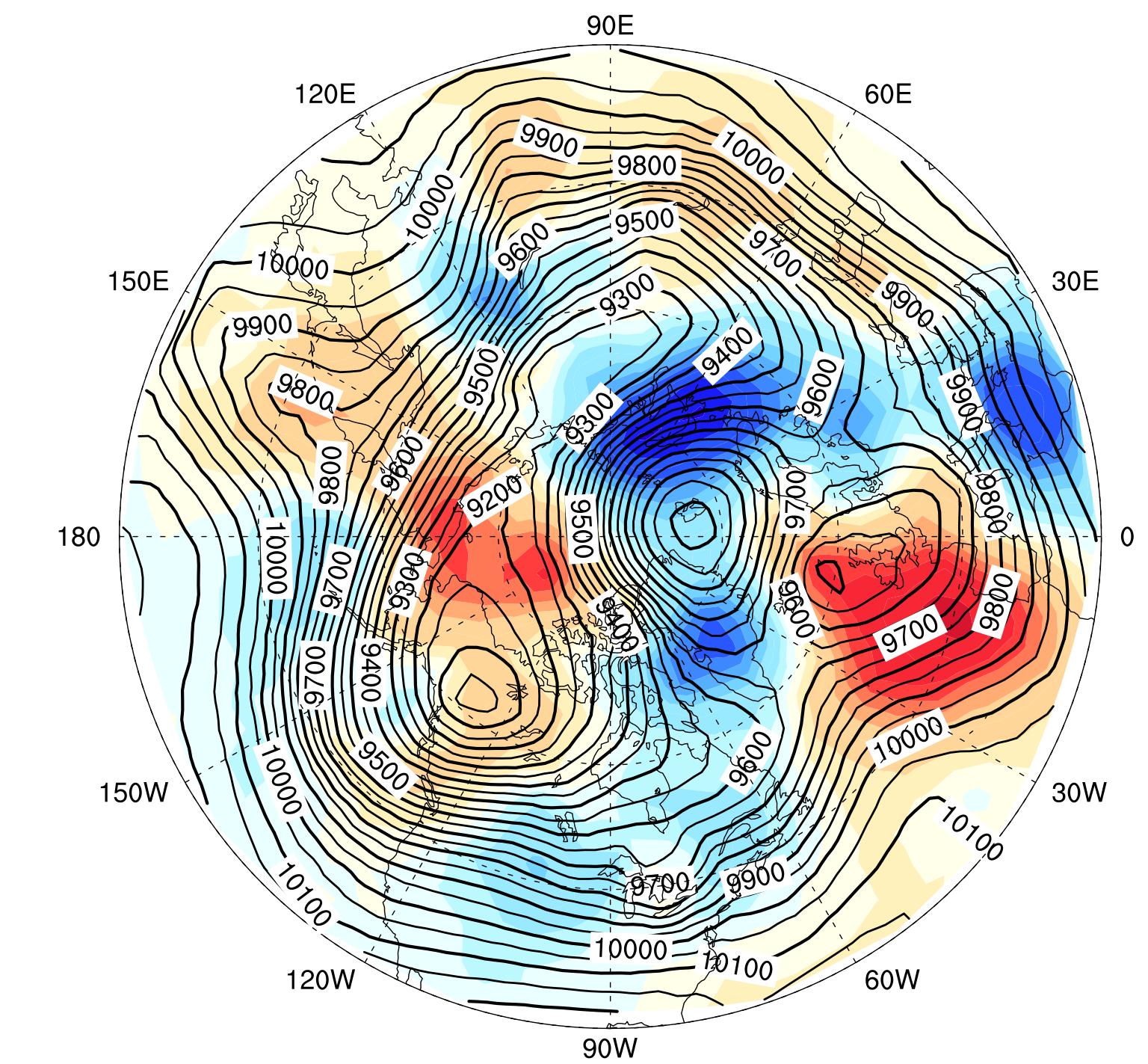
ME



SH

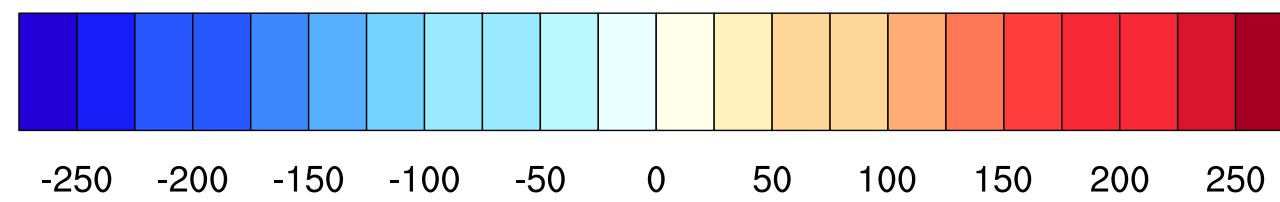


NI

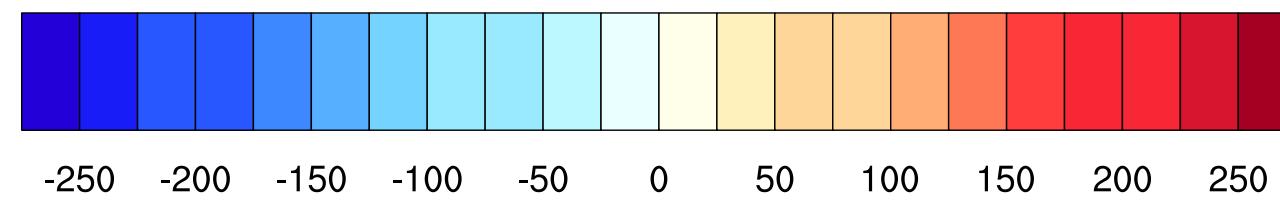


-250

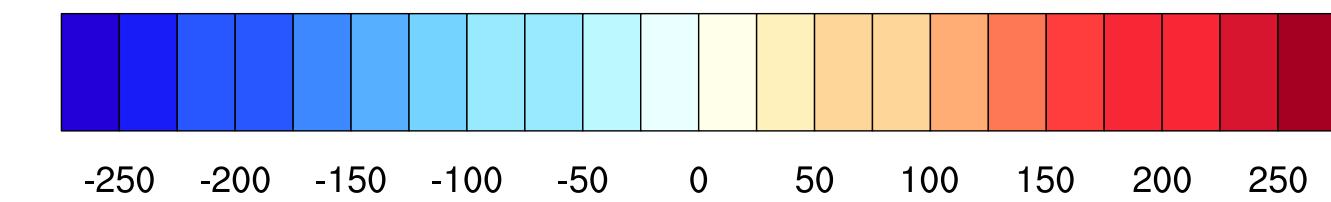
250



CONTOUR FROM 9050 TO 10250 BY 50



CONTOUR FROM 9050 TO 10250 BY 50

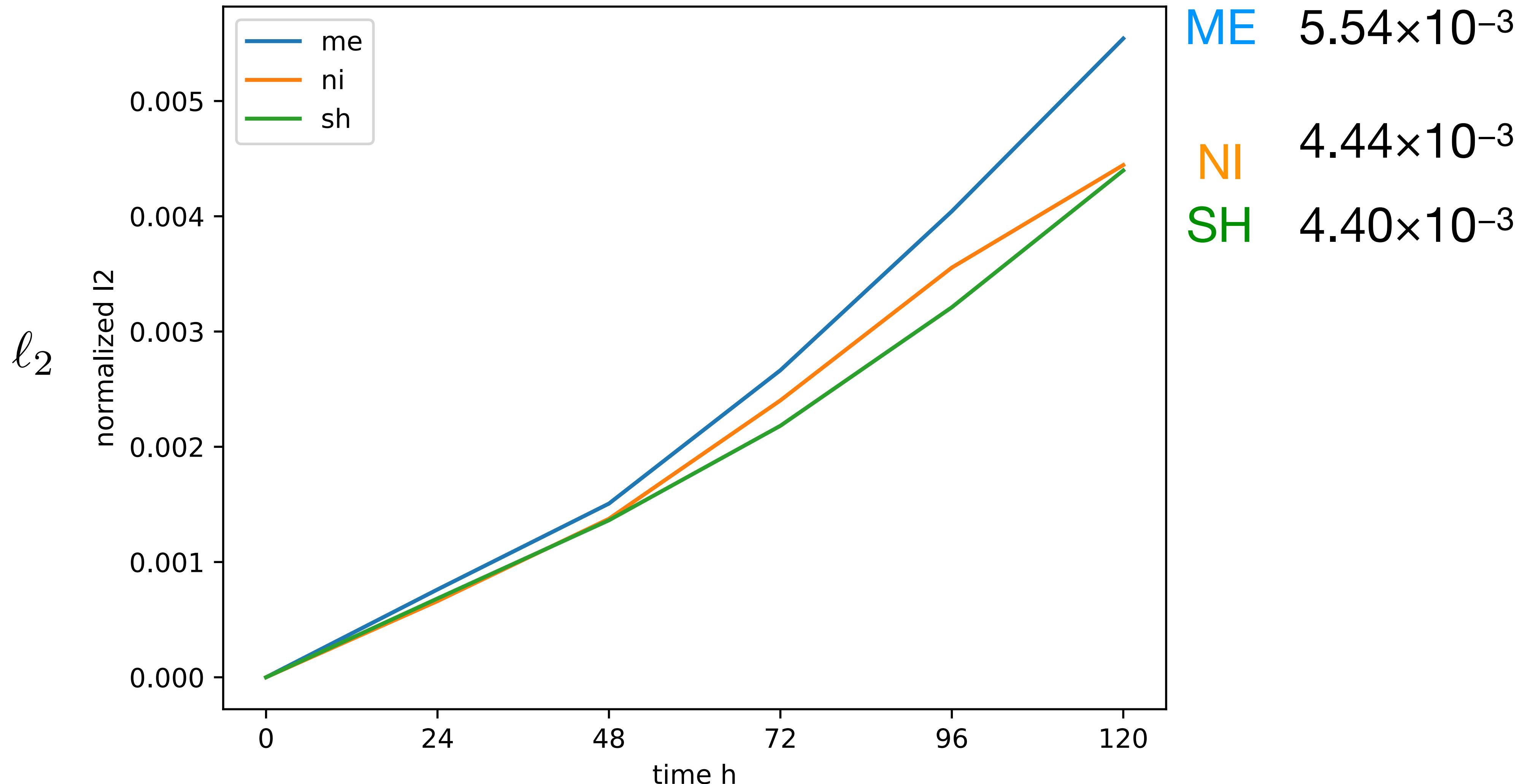


CONTOUR FROM 9050 TO 10250 BY 50

# Case 7: Forecast from 500 hPa analysis

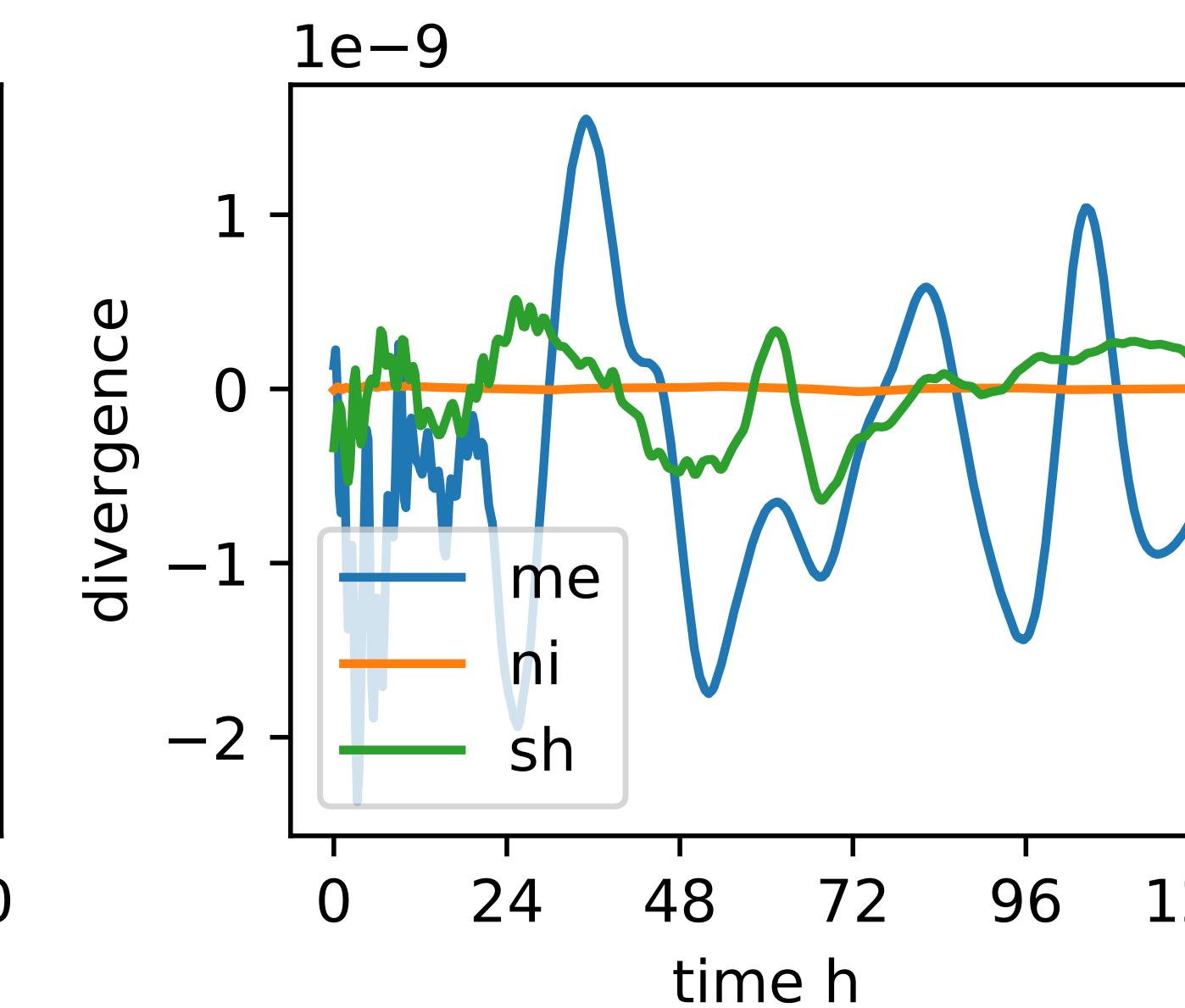
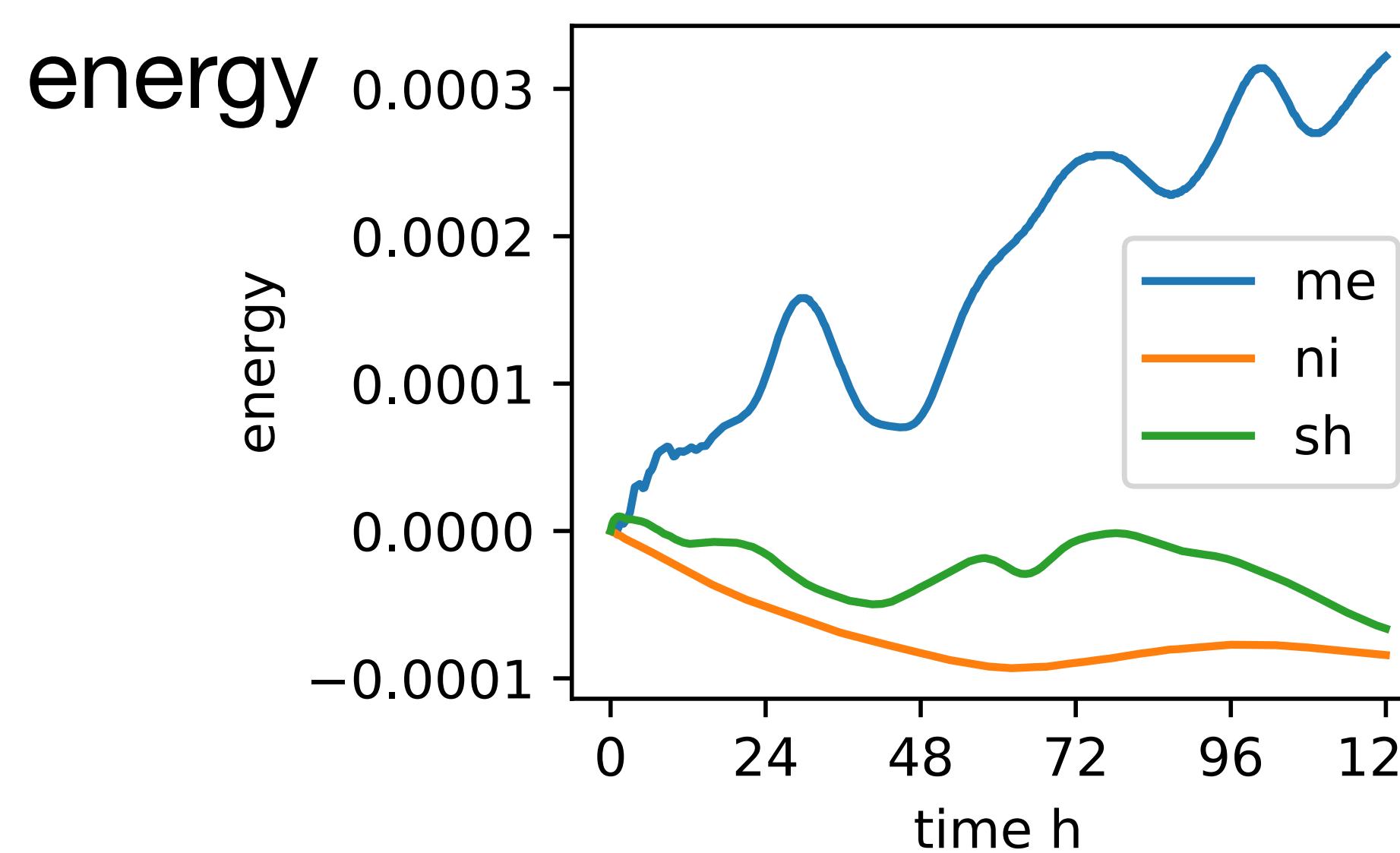
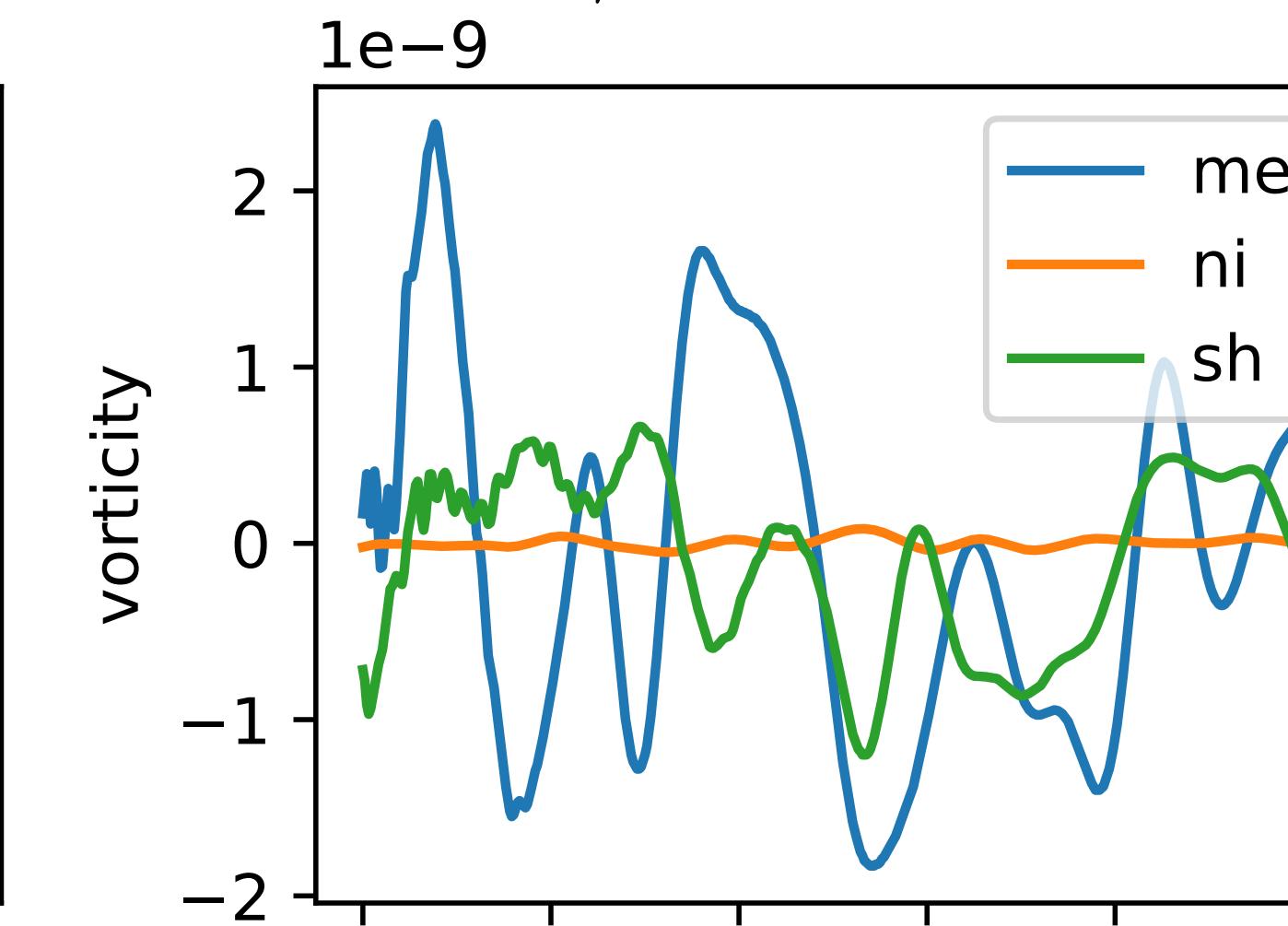
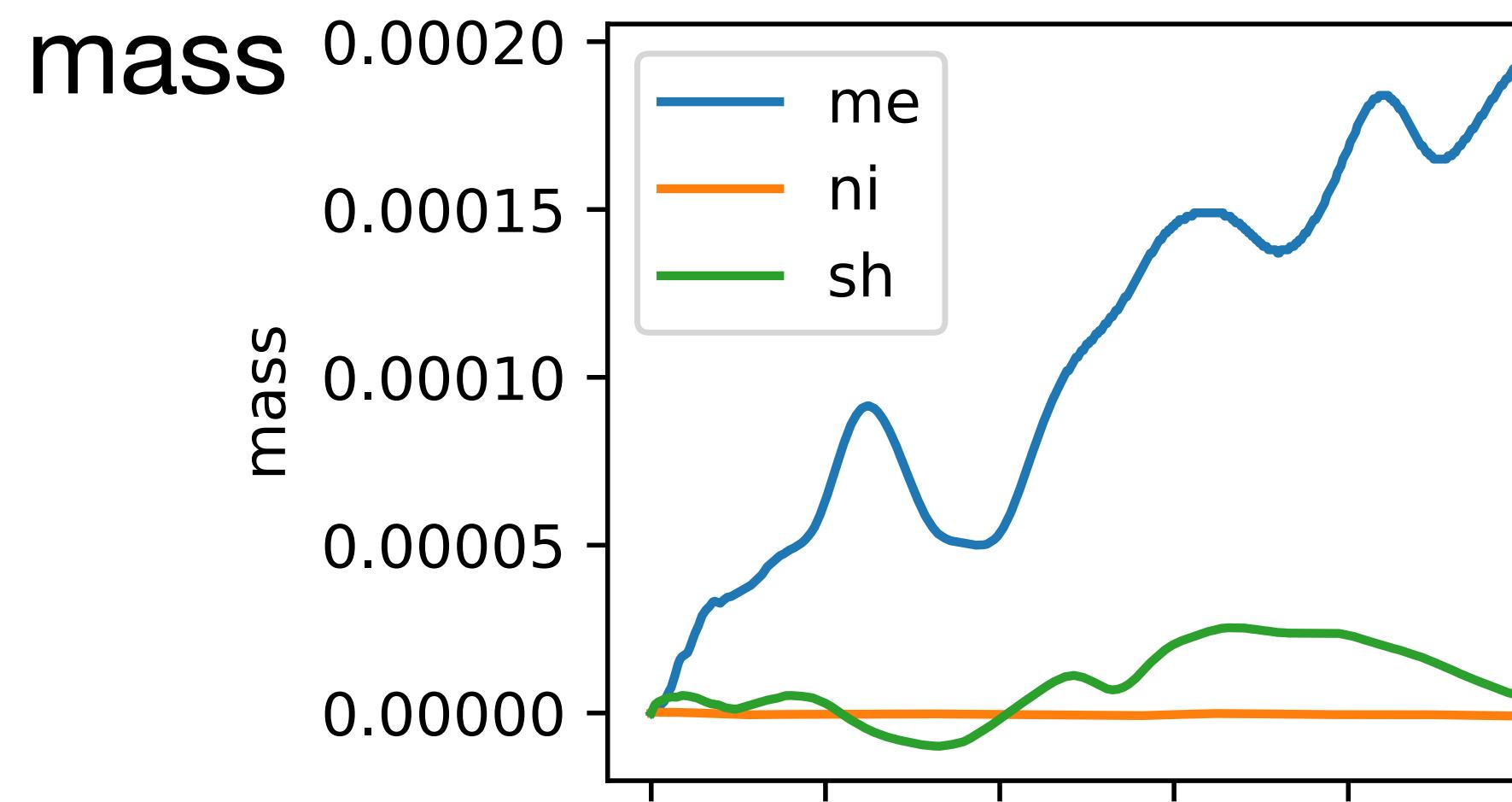
$$n = 2562 \Delta t = 15 \text{ min } \varepsilon = 4.25 \gamma = -2 \times 10^{-7}$$

h error from T213



# Case 7: Forecast from 500 hPa analysis

$$n = 2562 \Delta t = 15 \text{ min } \varepsilon = 4.25 \gamma = -2 \times 10^{-7}$$



# Settings and l2 norms

case	2	3	5	6	7
$\Delta t$ m		24			15
$\varepsilon$		4.75	4		4.25
$\gamma$		NA	-2E-08		-2E-7
NI	2.15E-09	3.35E-09	7.67E-04	6.05E-03	4.44E-03
ME	5.39E-09	1.62E-08	7.89E-04	6.81E-03	5.54E-03
SH	7.91E-10	6.79E-09	7.82E-04	8.01E-03	4.40E-03

# Summary

- Minimum Energy: “Scars”, Spherical Helix: Poles, Icosahedron: Vertices
- In steady state tests, RBF is not as robust as the spectral method but very accurate.
- Spherical helix and NICAM nodes are better than minimum energy nodes, but the differences are less significant in realistic tests.
- NICAM nodes have a good conservation of vorticity and divergence.

# Q&A

Q. Subich: What are possible causes of the differences in time evolution of error in conservation?

A. It may be related to the inhomogeneity such as scars in ME.

Q. Behrens: Have you tested other RBF?

A. GA seems to be more sensitive to the shape parameter than ME.

Q. Côte: What is the computational complexity?

A.  $O(n^3)$  to solve linear systems in precomputation and  $O(n^2)$  matrix operations per time step. Techniques exist to reduce computational loads.